# On emission measures and element densities and masses inferred from XSPEC

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# RESULTS

## INTRODUCTION

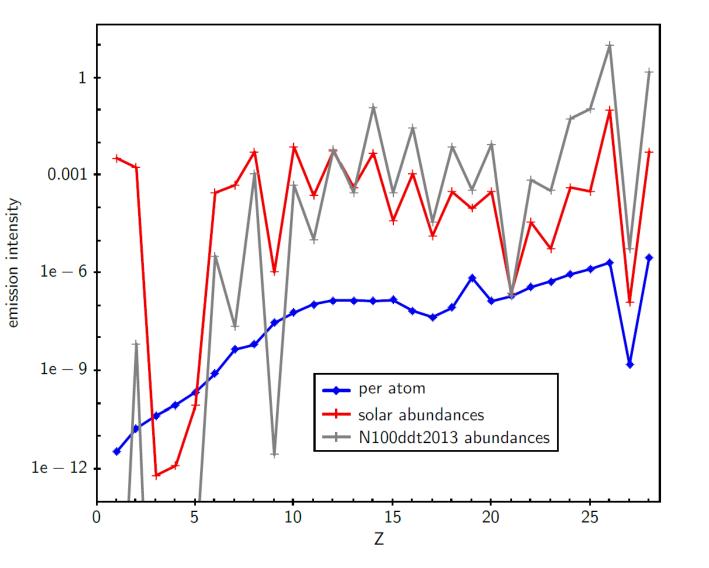
Spectral model fits to observed spectra of hot (X-ray emitting) plasmas provide us crucial information on the state and amount of hot plasma. The spectral fits inform us of the elemental composition of the plasma and its temperature and, in the case of non-equilibrium ionization plasmas, tell us the age of the plasma by the parameter  $\tau$  = ne t, with t the age since the plasma was shock heated.

The most common modeling software for X-ray spectra is the XSPEC software package (Arnaud 1996). The purpose of the current work is to explain the assumptions and provide corrections for cases where the assumptions are not accurate.

Definition of emission measure (EM) and definition of abundance factors (A\_Z) ) given below, with A\_sol a solar abundance set:

> $EM_e = \int n_e(\vec{r})^2 dV = n_e^2 V \quad \text{(if uniform)}$  $EM_H = \int n_e(\vec{r}) n_H(\vec{r}) dV = n_e \ n_H \ V \ \text{(if uniform)}$  $EM_Z = \int n_e(\vec{r}) n_Z(\vec{r}) dV = n_e \ n_Z \ V \ \text{(if uniform)}$

First, we review the contributions of various elements to the observed X-ray spectrum. In general, elements that strongly contribute (like Fe and Ni) are tightly constrained, but weakly contributing elements like H, He, Li, Be, B, C are poorly constrained, as shown in Fig,1 below.



#### **Table 4.** Fits to simulated spectra with added H and $He^{a}$

Spectrum	$A_H, A_{He}$	$\chi^2$	probability	$A_H, A_{He}$	$\chi^2$	probability
LSN	2.6E-16,7.5E-05	110.4	0.167	1, 1	109.8	0.177
	$^{4,4}$	111.2	0.153	8,8	114.4	0.109
	16,16	119.2	0.063	24,24	122.4	0.042
	32,32	124.6	0.031	40,40	122.4	0.025
MSN	2.6E-16, 7.5E-05	90.6	0.662	1, 1	105.6	0.258
	1.4, 1.4	113.1	0.127	1.8, 1.8	119.2	0.063
	2.2, 2.2	124.7	0.033	2.6, 2.6	128.5	0.018
HSN	2.6E-16,7.5E-05	93.2	0.590	2.6E-16,1	106.8	0.233
	1,7.5E-05	142.4	1.9E-03	1,1	176.1	1.6E-06
	0.1,0.1	96.0	0.509	0.2, 0.2	100.3	0.390
	0.3,0.3	106.7	0.236	0.4, 0.4	114.4	0.109
	0.5,0.5	123.0	0.038			

 $^{a}$ LSN has exposure 10ks, norm 2.32  $10^{-3}$ ; MSN has exposure 20ks, norm 6.96  $10^{-3}$ ; HSN has exposure 20ks, norm 2.32  $10^{-2}$ . All have kT=2.6 keV,  $n_e t = 1.31 \ 10^{12}$  and N100 abundances from column 8 of Table 2. H and He are added to the fit model, as listed.

0.01

 $n_H = A_1 n_{H,0}$  $n_Z = A_{sol,Z} A_Z n_{H,0}$ 

For XSPEC, the value of the electron density,  $n_{e,X}$ , is tied to the hydrogen density,  $n_{H,X}$ , by:

 $n_{e,X} = 1.2 \ n_{H,X}$ 

### ANALYSIS

The variation with solar abundance set is small (~3.5%). However the variation of the electron density between the different abundance sets (A\_Z) is large, with a number of examples given in Table 1. The ionization state and composition gives the ratio ne/nH,0. The variation is nearly a factor of 1000, instead of the single value assumed by XSPEC.

<b>Table 1.</b> Electron-to-(fiducial)hydrogen density ratio $n_e/n_{H,0}$ for different cases of									
composition and different temperatures.									
kT(keV):	0.10	0.25	0.50	1.0	2.0	4.0	10.0	40.0	
COMPOSITION									
SOLAR CASES <sup><math>a</math></sup>									
angr	1.20532	1.20652	1.20769	1.20836	1.20866	1.20880	1.20885	1.20886	
aspl	1.17672	1.17756	1.17829	1.17871	1.17893	1.17903	1.17907	1.17908	
feld	1.20536	1.20647	1.20765	1.20829	1.20856	1.20867	1.20872	1.20873	
aneb	1.16981	1.17092	1.17198	1.17258	1.17285	1.17297	1.17302	1.17303	
grsa	1.17861	1.17962	1.18058	1.18113	1.18138	1.18149	1.18154	1.18155	
wilm	1.20152	1.20224	1.20295	1.20335	1.20354	1.20362	1.20365	1.20366	
lodd	1.16470	1.16552	1.16625	1.16666	1.16686	1.16696	1.16670	1.16670	
lpgp	1.17506	1.17589	1.17668	1.17713	1.17736	1.17746	1.17775	1.17775	
lpgs	1.20159	1.20254	1.20342	1.20394	1.20420	1.20431	1.20436	1.20437	
NON-SOLAR CASES									
H $0^b$	0.20532	0.20652	0.20769	0.20836	0.20866	0.20880	0.20885	0.20886	
(H,He) $0.1^b$	0.12946	0.13066	0.13183	0.13250	0.13280	0.13294	0.13299	0.13300	

Figure 1. Illustration of the relative contributions of different elements to the 0.3 to 8 keV energy range for a spectrum, using the response matrices for the AstroSat SXT instrument (Singh et al. 2017). The maximum intensity in photons s<sup>-1</sup> keV<sup>-1</sup> is shown for elements Z = 1 to 28 for three cases: per atom, for solar abundances and for N100ddt2013 abundances. For each case the vertical scale is arbitrary. The XSPEC vvapec model was used with temperature of 1 keV. For N100ddt2013, H, Li, Be and B not shown because they are so small: they at  $2.9 \times 10^{-22}$ ,  $1.9 \times 10^{-24}$ ,  $4.1 \times 10^{-33}$ , and  $1.2 \times 10^{-16}$  respectively times the contribution of Fe.

Because the abundances of H and He are normally very large, they can contribute significantly to the X-rays spectrum. This is illustrated in Fig.2 below, i.e. H abundance needs to be 30 times solar for H to contribute equally to Fe. Because solar abundance of some elements is very low (e.g. Li) the abundance relative to solar to contribute equally as Fe is large (~10^11).

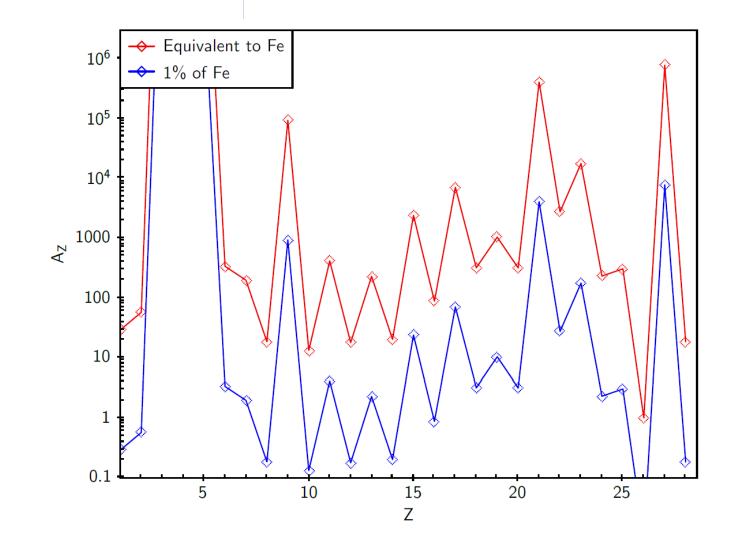


Figure 2. The abundance factors  $A_Z$  required to yield the same maximum contribution (red line) or to yield a contribution 1% as large (blue line) as the contribution from Fe to the 0.3 to 8 keV X-ray spectrum. The response matrices for the AstroSat SXT instrument was used and an equilibrium spectrum with temperature kT=2 keV and the abundance factor of Fe was set to 1 for the same-contribution line (or 0.01 for the 1%) line). The abundance factors for Li, Be and B off scale, at values of  $1.5 \times 10^{11}$ ,  $7.8 \times 10^{10}$ , and  $1.1 \times 10^{9}$  for equal contribution, with values 0.01 times as large for a 1% contribution.

A number of test cases with prespecified abundances were carried out to illustrate the changes in density that results for non-solar abundance case, which are caused by the correction from including the correct values of electron density.

Figures 4 and 6 compare cases with solar abundance to those with light elements removed. With e.g H and He removed the element densities increase by factors of 10 and masses increase by the same factor.

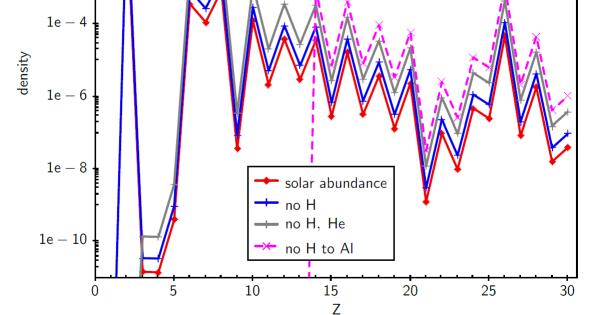


Figure 4. Element densities for different cases. The four cases shown are solar abundance, no H, no H and He, and no H through Al.

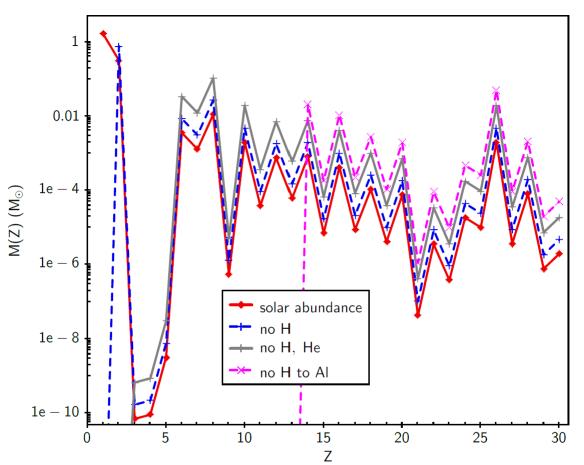
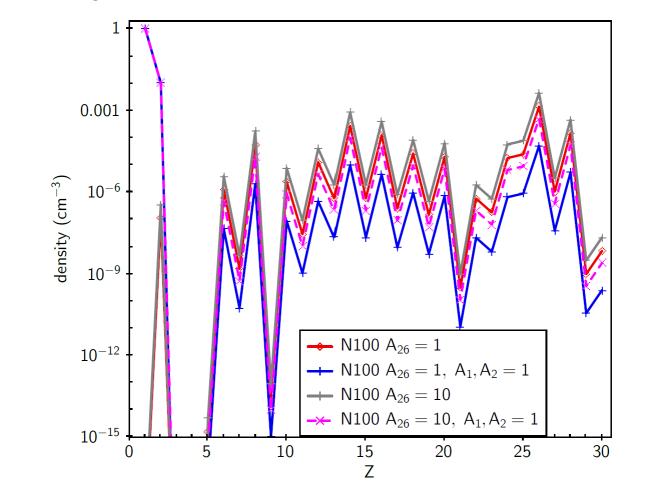


Figure 6. Masses for different cases. For calculation, the distance is 1 kpc, the emitting volume is a sphere of radius 2 pc, and the XSPEC norm is 0.1. The four cases shown are solar abundance, no H, no H and He and no H through Al.



Figures 5 and 7 compare cases

$N100Fe1HHe^{d}$	1.1965	1.1970	1.1971	1.1973	1.1974	1.1975	1.1975	1.1976
$N100 Fe 10^{e}$	0.01080	0.01574	0.01720	0.01862	0.02021	0.02112	0.02150	0.02152
$\rm N100Fe10HHe^{\it f}$	1.2062	1.2111	1.2126	1.2140	1.2156	1.2165	1.2169	1.2169

 $0.00992 \quad 0.01112 \quad 0.01229 \quad 0.01296 \quad 0.01326 \quad 0.01340 \quad 0.01345 \quad 0.01346$ 

0.00743 0.00833 0.00935 0.01000 0.01030 0.01043 0.01049 0.01050

 $0.00093 \quad 0.00142 \quad 0.00155 \quad 0.00168 \quad 0.00183 \quad 0.00192 \quad 0.00195 \quad 0.00196$ 

 $0.00108 \quad 0.00157 \quad 0.00172 \quad 0.00186 \quad 0.00202 \quad 0.00211 \quad 0.00215 \quad 0.00215$ 

#### <sup>a</sup>Abundance references given in the XSPEC manual- see the abund command.

<sup>b</sup> All other elements taken to have solar abundance factors,  $A_Z = 1$ .

<sup>c</sup> N100 model (Seitenzahl et al. 2013) with  $A_{Fe}=1$ .

 $^{d}$ N100 model with  $A_{Fe}=1$ , H He added at solar abundance.

 $^{e}$  N100 model with  $A_{Fe}=10$ .

 $(H,He) 0^b$ 

(H to N)  $0^b$ 

(H to Al)  $0^b$ 

 $N100Fe1^{c}$ 

 $^{f}$ N100 model with  $A_{Fe}$ =10, H He added at solar abundance.

To simplify notation, hereafter we define the electron density ratio as  $r_e$ , given by:

$$r_e = \frac{n_e}{1.2 n_{H,0}}$$
 = ratio of electron density to XSPEC value

the fiducial density is given in terms of  $norm_X$  by:

$$n_{H,0} = r_e^{-1/2} \sqrt{10^{14} \ norm_X \ \frac{4 \ \pi \ D_A^2}{1.2 \ V}}$$

The XSPEC norm for hot plasma emission models is proportional to  $EM_{H,X}$ :

$$norm_X = 10^{-14} EM_{H,X} / (4\pi [D_A(1+z)]^2)$$
 with  
 $EM_{H,X} = \int n_{e,X}(\vec{r}) n_{H,X}(\vec{r}) dV = n_{e,X} n_{H,X} V$  (if uniform

From the detailed calculation (Leahy, Foster, Seitenzahl, 2024) we show the correct density of the plasma is given by:

$$n_{H} = A_{1} r_{e}^{-1/2} n_{H,X}$$
$$n_{Z} = A_{Z} A_{sol,Z} r_{e}^{-1/2} n_{H,X}$$

Element masses depend on the element densities and emitting volume,

 $M_Z = m_Z \int n_Z(\vec{r}) dV = m_Z n_Z V \text{ (if uniform)}$  $= m_Z \ EM_Z/n_e = m_Z \ A_{sol,Z} \ A_Z \ EM_{H,X}/n_e$  $= m_Z A_{sol,Z} A_Z \sqrt{\frac{10^{14} norm_X 4 \pi D_A^2 V}{10^{14} norm_X 4 \pi D_A^2 V}}$ 

We carried out a number of simulations in XSPEC simulating observed spectra using the Astrosat SXT response matrices. These were carried out for typical SNR fluxes, based on known type la SNRs with observed X-ray spectra, shown in Table

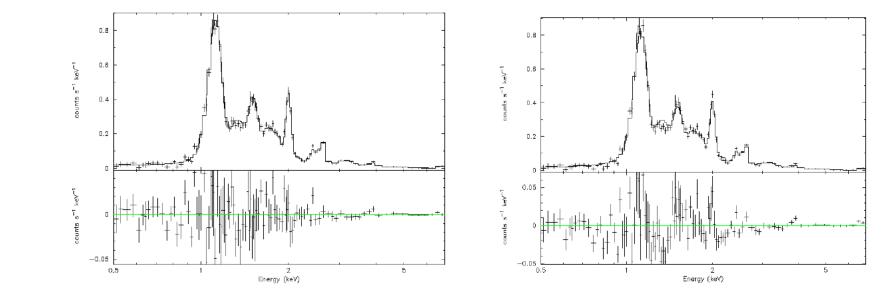
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#### **Table 3.** Measured temperatures and EM for type Ia SNRs<sup>a</sup>

SNR ID	type	$\mathrm{D}(\mathrm{kpc})$	$EM(10^{58}{ m cm}^{-3})$	norm	kT(keV)	$Age(yr)^b$	$density(cm^{-3})^b$	$n_e t (\mathrm{cm}^{-3}\mathrm{s})^b$
G53.6-2.2	Ia	7.8	0.022	3.02E-04	3.9	44600	0.096	5.41E + 11
G299.2-2.9	Ia	5	0.029	9.70E-04	1.36	8800	0.022	2.45E + 10
G306.3-0.9	Ia	20	1.8	3.76E-03	1.51	12800	2.5	4.04E + 12
G315.4-2.3	Ia	2.8	0.046	4.91E-03	3.04	11600	0.247	3.62E + 11
G352.7-0.1	Ia?	7.5	0.11	1.64E-03	3.2	7600	1.66	$1.59E{+}12$
average:				2.32E-03	2.602			1.31E + 12

<sup>a</sup>Values are from Tables 1 and 2 of Leahy et al. (2020), with references for EM, norm and kT given in Table 1. <sup>b</sup> Values for age and density are from Table 2 of Leahy et al. (2020), with  $n_e t$  calculated using post-shock density of  $4n_e$ .

2 of the simulated X-ray spectra and their spectral fits are shown in Fig. 3. The results from a set of simulated spectra and fits to those spectra are shown in Table 4. These verify that the abundances A\_Z derived from spectral fitting are quite uncertain, and the spectra are fit with significant changes in the H and He abundances.



with realistic type Ia abundances from the N100 model. The difference is that the spectral fit is normalized with Fe abundance =1 (red line) or Fe abundance = 10(grey line). These two spectral fits are compared to the same abundance set but with added H and He at solar abundance, which is orders of magnitude large that that produced by the N100 model, but still comparable with what might be assumed for a spectral fit with X\$PEC. The densities and masses for added H and He are lower by factors of ~30 to 60 than for the case using realistic H and He abundances. I.e. assuming too high H and He results in significant underestimate of densities and abundances, compared to the real physical case.

Figure 5. Element densities for N100 model (Seitenzahl et al. 2013), with and without added H and He. For calculation, the distance is 1 kpc, the emitting volume is a sphere of radius 2 pc, and the XSPEC norm is 0.1, which yields XSPEC  $EM_X = 1.2 \times 10^{57}$  cm<sup>-3</sup>. The four cases shown are N100 abundances normalized with  $A_{Fe}=1$ , N100 with  $A_{Fe}=1$  and H and He added at solar abundances, N100 abundances normalized with  $A_{Fe}=10$ , and N100 with  $A_{Fe}=10$  and H and He added at solar abundances.

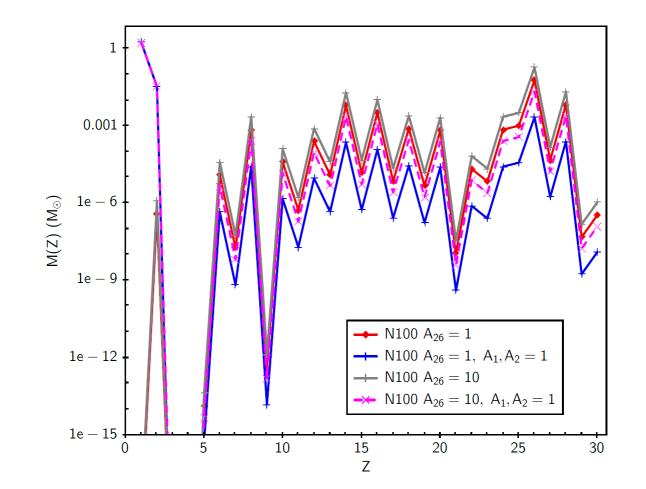


Figure 7. Masses for different cases. For calculation, the distance is 1 kpc, the emitting volume is a sphere of radius 2 pc, and the XSPEC norm is 0.1. The four cases shown are N100 abundances normalized with  $A_{Fe}=1$ , N100 with  $A_{Fe}=1$  and H and He added at solar abundances, N100 abundances normalized with  $A_{Fe}=10$ , and N100 with  $A_{Fe}=10$  and H and He added at solar abundances.

## **CONCLUSIONS/FUTURE WORK**

For cases of non-solar abundances, with A1 or A2 different from 1, significant corrections are needed to the element densities to account for the difference between ne and ne,X . ne/ne,X can be much less than 1 as illustrated by the examples in Table 1. We define the electron density ratio re. The corrected densities are higher, by the factor re<sup>^</sup>-1/2, and the inferred element masses are corrected by the same factor (re<sup>^</sup>-1/2) thus can also be much larger. This implies that masses that may have been derived in the past using norm's from XSPEC spectral fitting on hydrogen poor plasmas, such as expected for Type Ia SNRs, may be significantly underestimated.

Figure 3. The HSN (high signal-to-noise, see text) simulated X-ray spectrum, which uses rmf, arf and background files for AstroSAT SXT, and abundances given by column 8 in Table 2 and model fits. The other spectrum parameters are  $N_H = 10^{21} \text{ cm}^{-2}$ , kT=2.6 keV,  $n_e t = 1.31 \times 10^{12} \text{ cm}^{-3}$  s and norm= $2.32 \times 10^{-2}$ . The fit model on the left has original N100 abundances, whereas the fit model on the right has added H and He with  $A_H = A_{He} = 1$ . Although the fit with solar H and He is statistically poor (probability=1.6E-06, Table 4), it is not clearly different except for larger residuals by a factor of  $\sim 1.2$ .