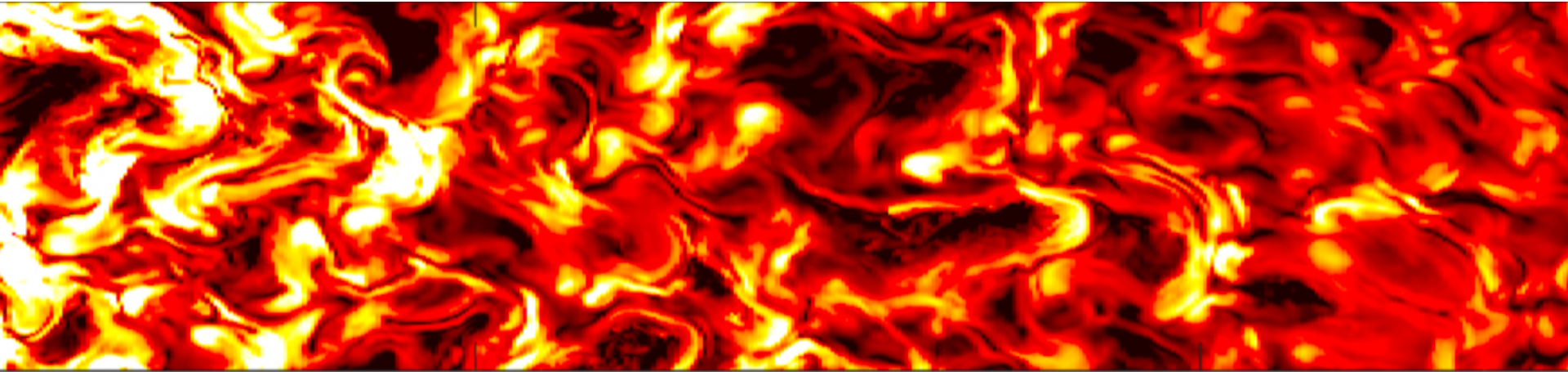


# Maximum energy cosmic rays from Galactic SNRs

simulations of quasi-parallel and -perpendicular shocks

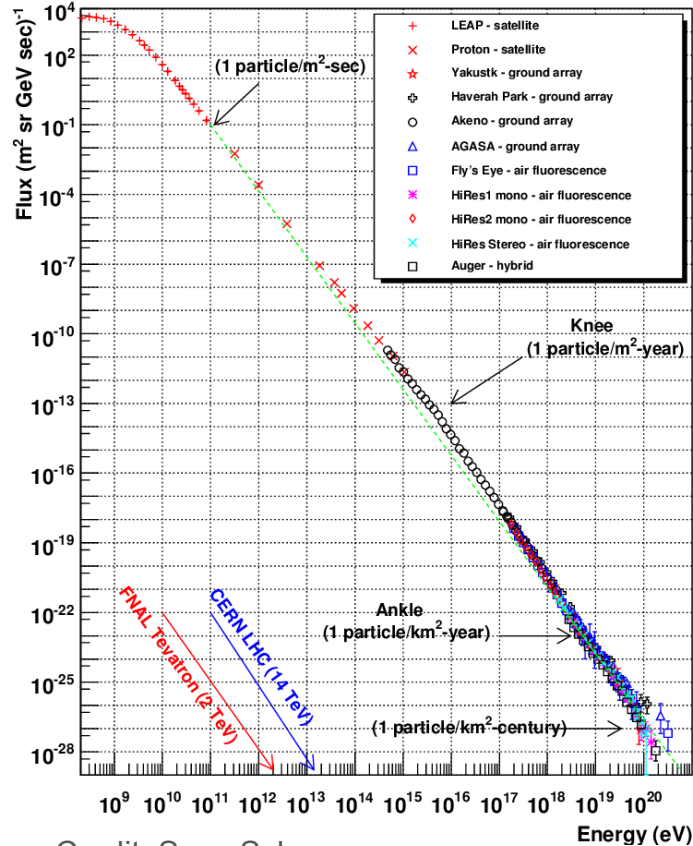


**Emily Simon**, University of Chicago  
Damiano Caprioli, University of Chicago  
Colby Haggerty, University of Hawaii, IfA  
Brian Reville, Max-Planck-Institut fuer Kernphysik

**SUPERNOVA REMNANTS III**  
AN ODYSSEY IN SPACE AFTER STELLAR DEATH  
9-15 June 2024, Chania, Crete, Greece



# The PeVatron Problem

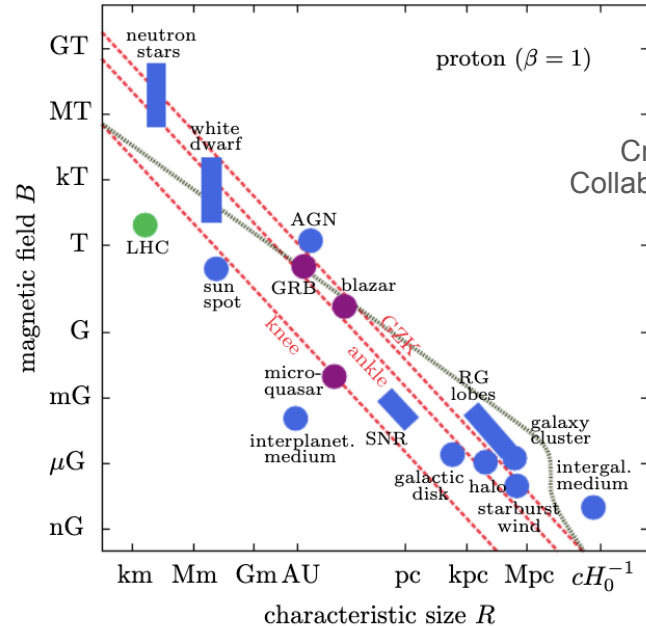


Credit: Sven Schoo

Region between 1 - 100 PeV not well understood

- Spectral break at  $\sim 4$  PeV
- Transition from Galactic to extragalactic

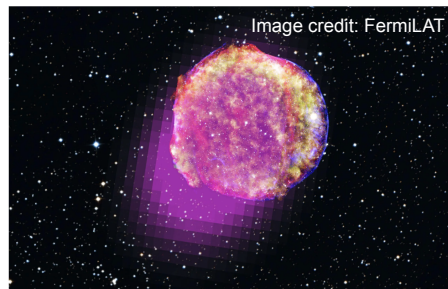
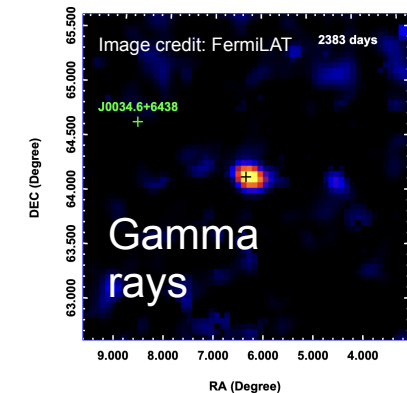
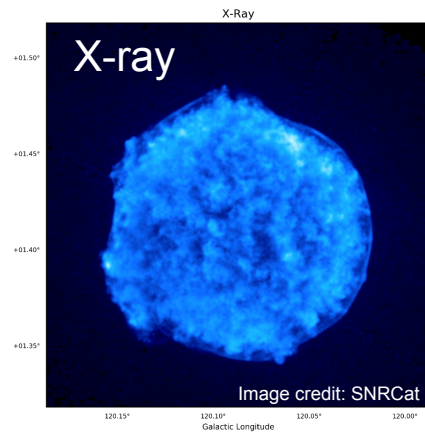
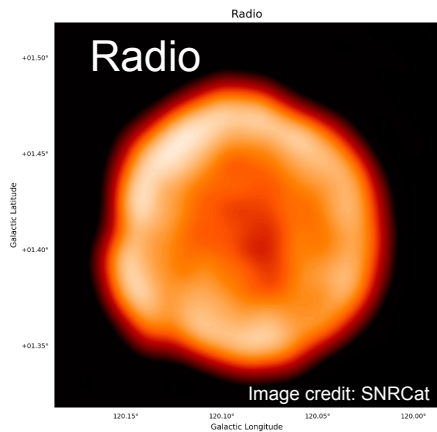
Origin of cosmic rays (CRs) in the knee region remains a mystery, but SNRs are well motivated



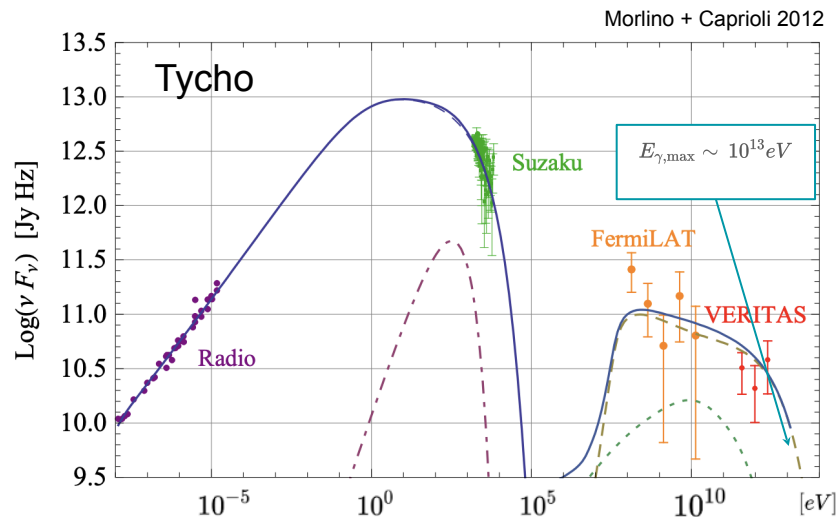
Credit: IceCube Collaboration, 2017

# The PeVatron Problem

G120.1+01.4



Gamma-rays detected by Fermi's LAT show that the remnant of Tycho's supernova shines in the highest-energy form of light. This portrait of the shattered star includes gamma rays (magenta), X-rays (yellow, green, and blue), infrared (red) and optical data.  
Credit: Gamma ray, NASA/DOE/Fermi LAT Collaboration; X-ray, NASA/CXC/SAO; Infrared, NASA/JPL-Caltech; Optical, MPA, Calar Alto, G. Krause et al. and DSS



$$p + p \rightarrow p + p + \pi^0$$

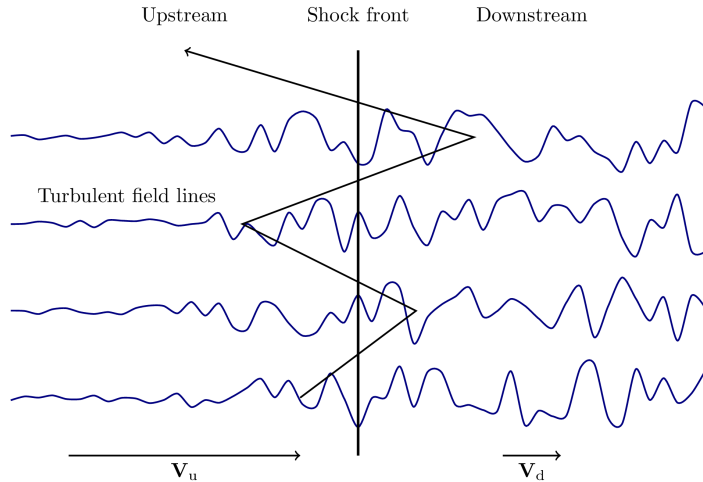
$$\pi^0 \rightarrow \gamma + \gamma$$

Each  $\gamma$  has  $\sim 10\%$  of the energy of its parent proton

This means the maximum energy proton had

$$E_p \sim 10^{14} eV$$

# Diffusive Shock Acceleration (DSA)



Credit: Marc Pulupa, Berkeley

*Lagage + Cesarsky, 1983:*

Bohm Diffusion in  $\sim 1 \mu G$  fields, saturate the resonant instability, and get CR energies up to hundreds of TeV

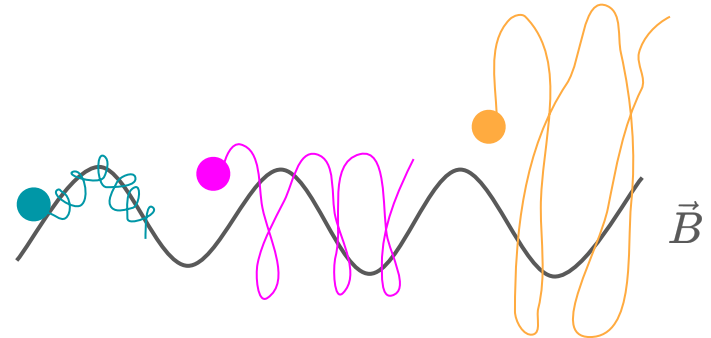
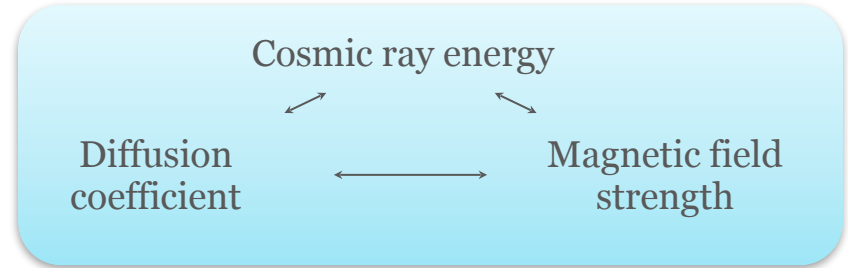
**Need magnetic field amplification!**

DSA in several hundred  $\mu G$  fields can achieve PeV CRs

The **maximum energy** is set by the rate at which particles come back to the shock

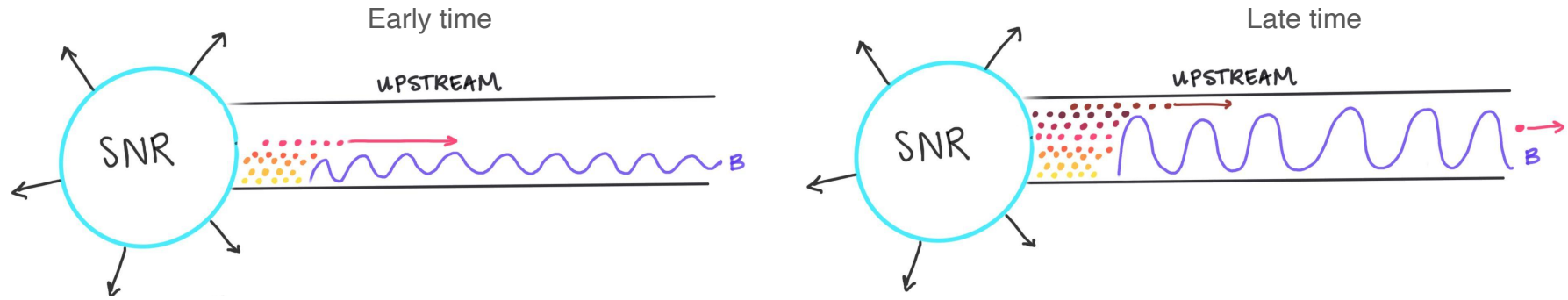
→ This depends on the strength of the magnetic fields

→ Which depends on the diffusion coefficient



# The Bell Instability

How do you make  $\sim 100 \mu G$  magnetic fields? *Bell 2004*:



Bell is the fastest instability, but how long does it take to make  $\sim 100 \mu G$  magnetic fields?

$$\text{Growth rate, } \gamma = \left( \frac{k j_{cr} B_0}{\rho} \right)^{1/2}$$

**Need  $\sim 5-10$  e-foldings** to achieve strong magnetic field growth

# The Bell Instability

*Bell et al. 2013:*

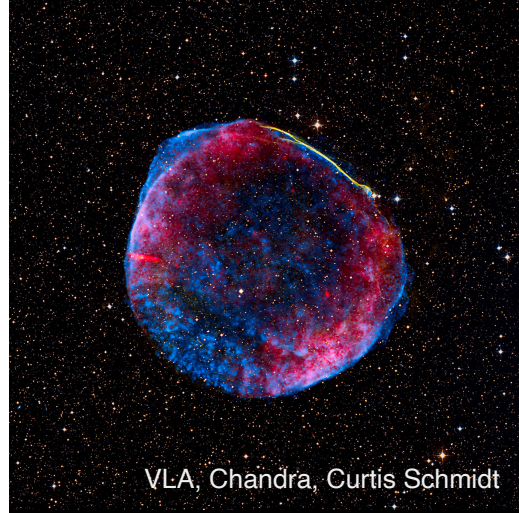
It takes ~days after the SNe to achieve 5-10 e-foldings of the Bell Instability, after which our magnetic field is strongly amplified

We can then calculate how the maximum energy grows in time in this amplified field and compare it to the lifetime of the SNR while the shock velocity is ~ constant

$$T_{max} = 0.005 \frac{P_{CR}}{\rho u_{sh}^2} \rho u_{sh}^3 t \sqrt{\frac{\mu_0}{\rho}} \approx 100 TeV$$

We want to test this from **first principles** to see if there's something missing

SN 1006  
~ 180 TeV

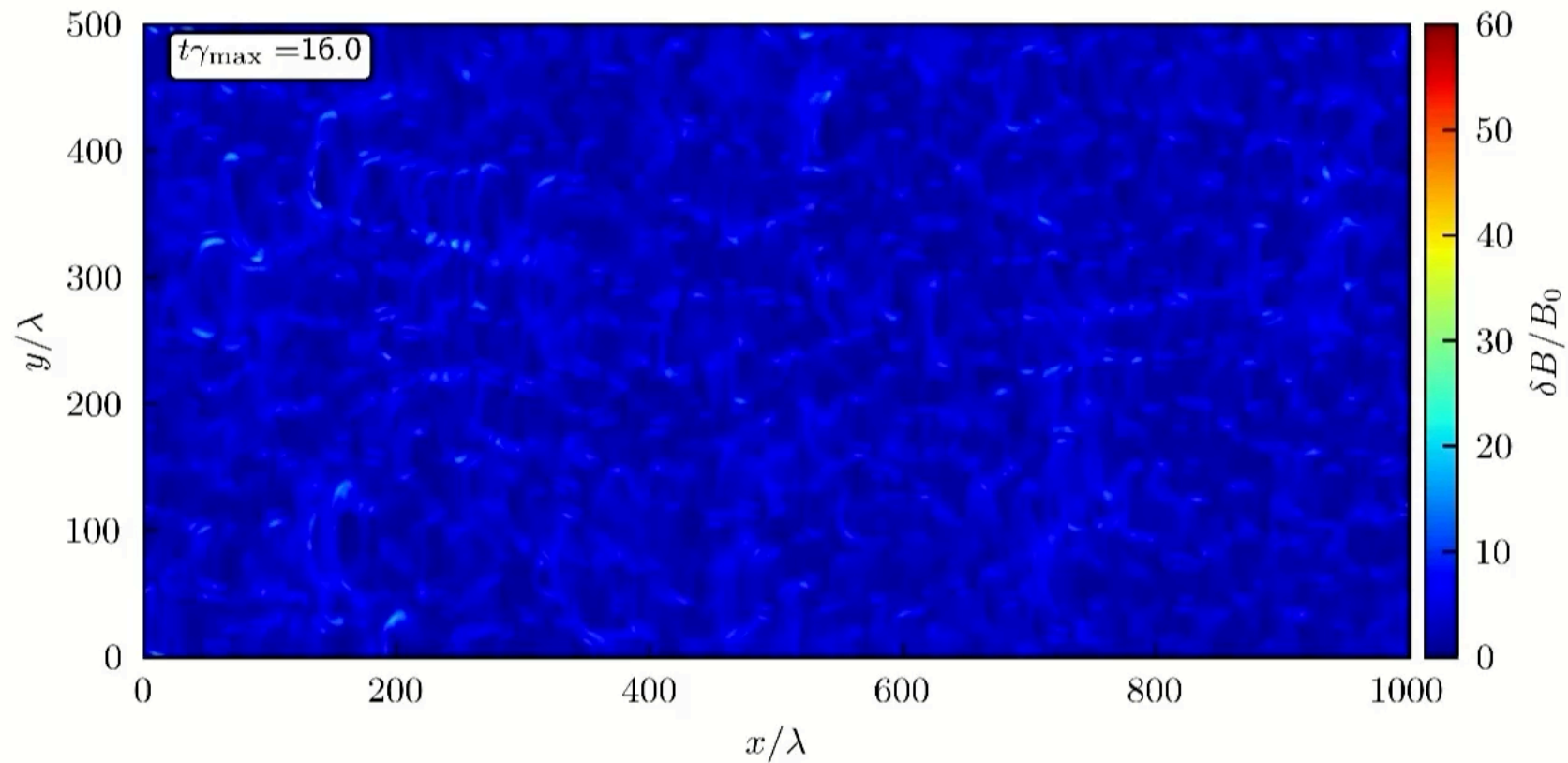


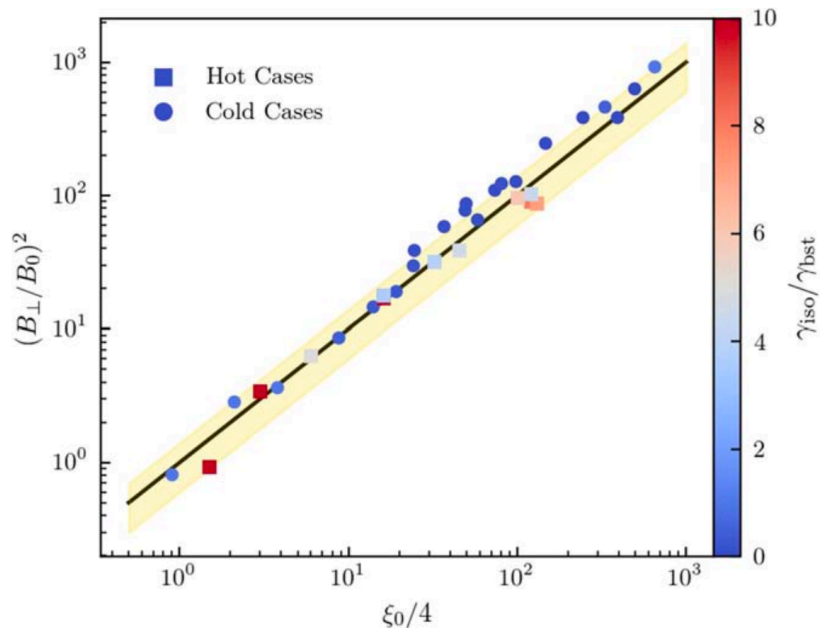
VLA, Chandra, Curtis Schmidt

Cas A  
~ 140 TeV



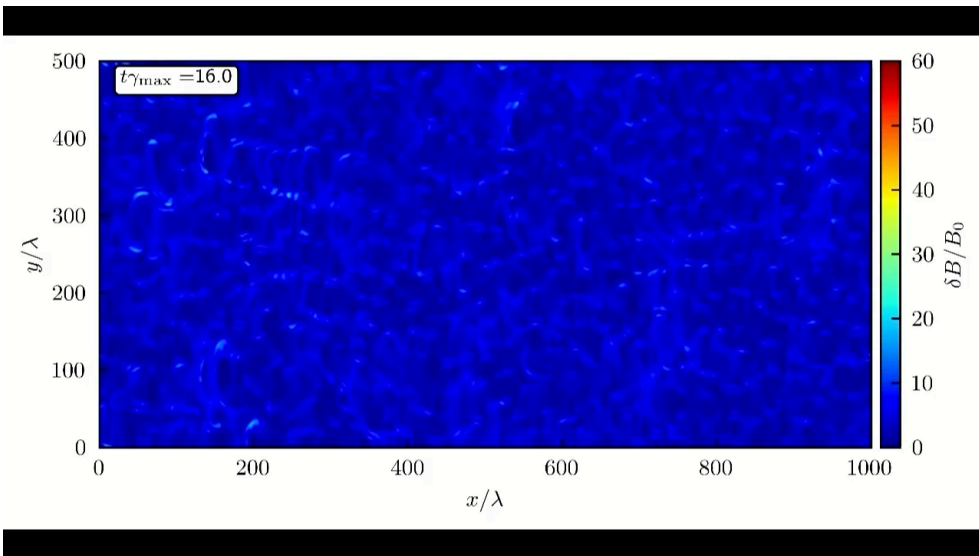
Chandra





$$\xi_0 = \frac{T^{01}}{P_{B,0}} \frac{v_{bst}}{c}$$

$$= 2\gamma_{iso}\gamma_{bst} \frac{n_{cr}}{n_g} \frac{v_{bst}^2}{v_{A,0}^2} \left( 1 + \frac{1}{3} \frac{v_{iso}^2}{c^2} \right)$$



Magnetic field saturation depends on the free momentum in CRs:

$$\xi_0 \approx 4 \left( \frac{\delta B}{B_0} \right)^2$$



# Simulations help us understand these processes

Particle In Cell (PIC): We use the code: **dHybridR** (Haggerty & Caprioli 2019)

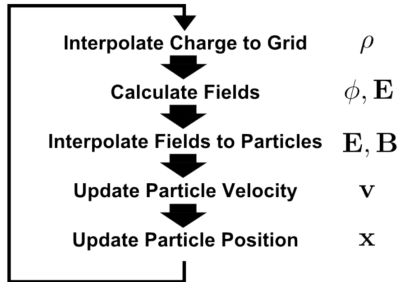
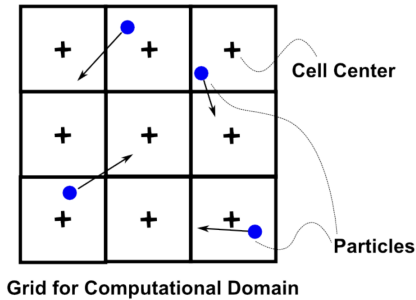
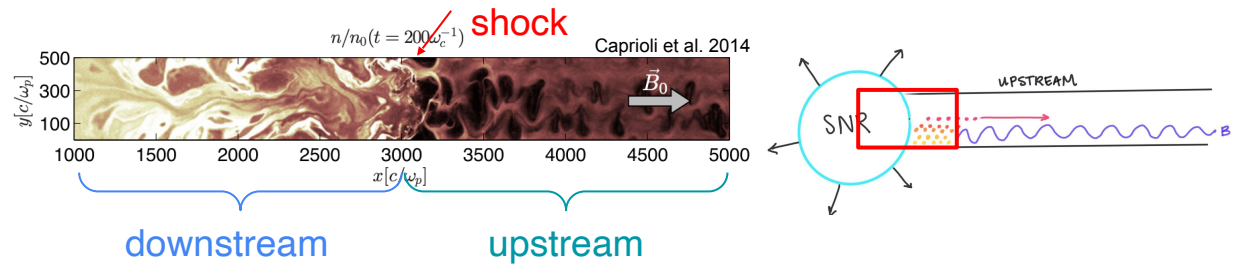
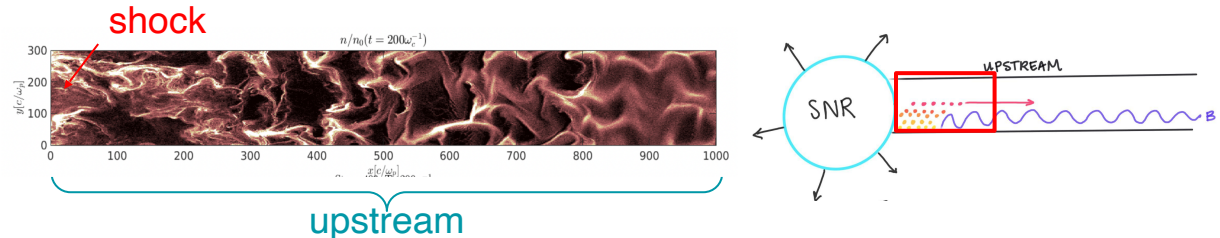


Image credit: Ebersohn et al. 2014

Traditional shock setup: reflecting boundary wall + colliding flows



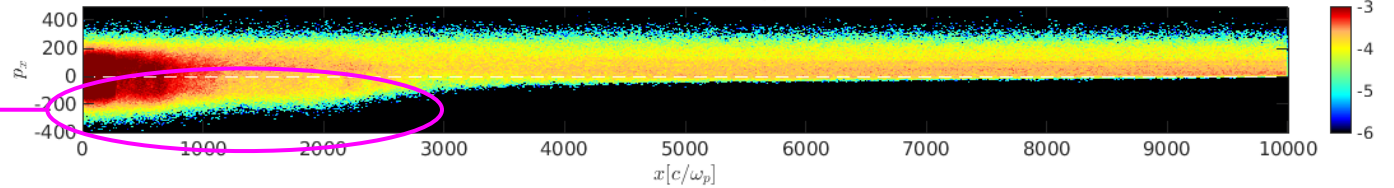
Novel “faux shock” setup: shock-like boundary wall



# Simulations Results

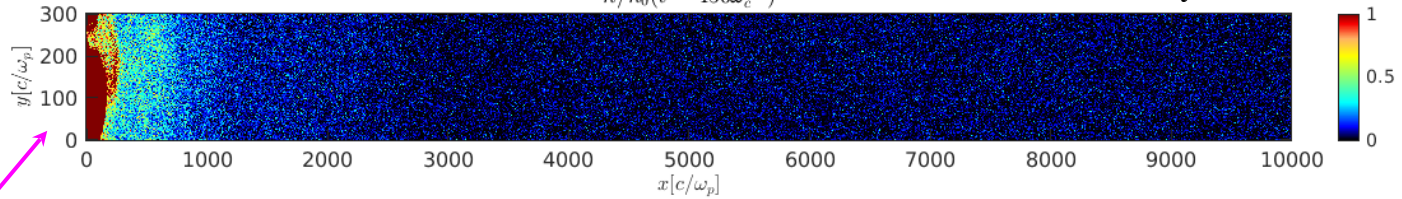
$\text{Log}[f(p_x)]t = 450\omega_c^{-1}$  x-directed momentum

CRs diffusing back towards the shock

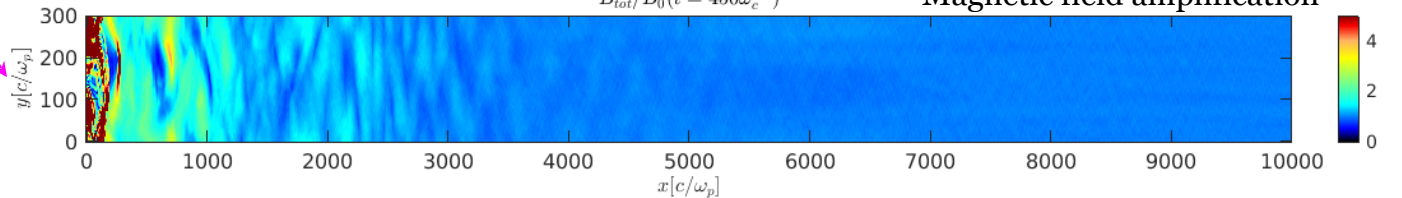


$n/n_0(t = 450\omega_c^{-1})$  CR density

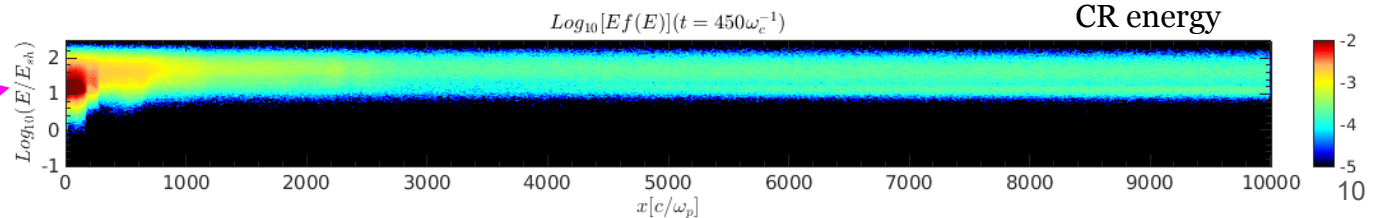
Greater CR density and stronger B field amplification closer to the "shock"



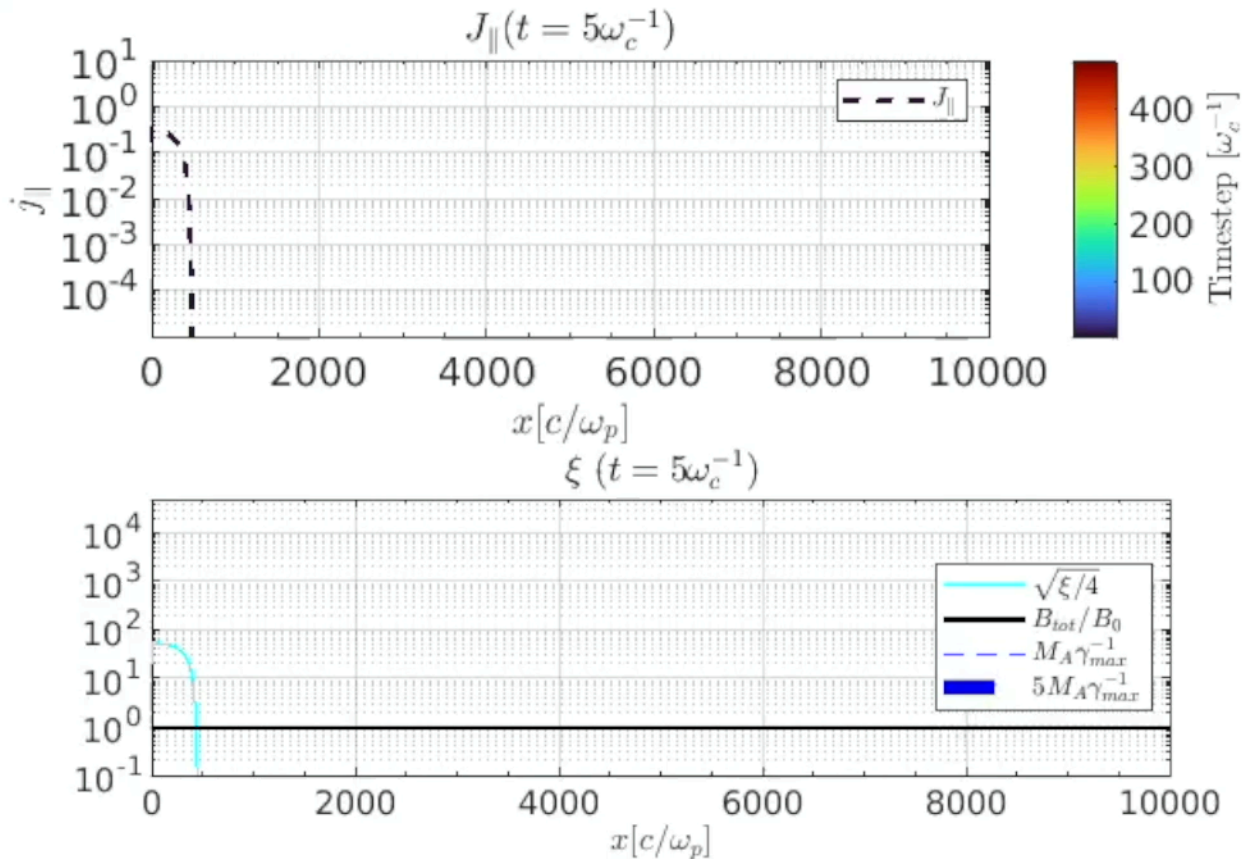
$B_{\text{tot}}/B_0(t = 450\omega_c^{-1})$  Magnetic field amplification



Highest energy CRs are in the far upstream and escaping the shock



# Current In Escaping Cosmic Rays



# Maximum Energy Cosmic Rays

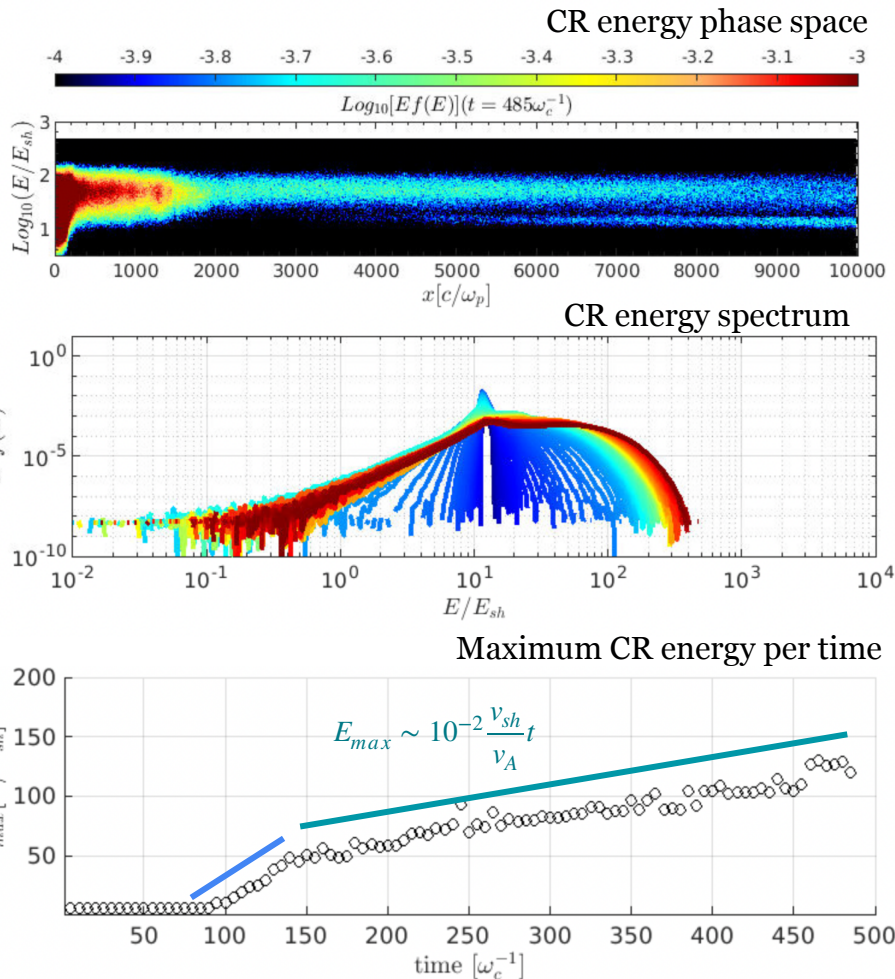
Which timescale is the bottleneck: acceleration or instability growth time?

$$\tau_{acc} \approx 6 \frac{D(E)}{v_{sh}^2} \approx 2 \frac{E_{max}}{m v_{sh}^2} \frac{1}{M_A^2} \frac{B_0}{B} \Omega_c^{-1}$$

$$\Delta t_{Bell} \approx 5 \gamma_{max}^{-1} = \left( \frac{1}{10} \frac{n_{cr}}{n_g} \frac{v_d}{v_{A,0}} \right)^{-1} \Omega_c^{-1}$$

$$\frac{\Delta t_{Bell}}{\tau_{acc}} \approx \frac{\frac{5\sqrt{3}}{2} \left(\frac{v_A}{c}\right)^{1/2} M_A^2}{\tilde{J}_{cr}^{1/2} E_{max}^{1/2}} \gtrsim 1$$

The instability growth timescale tends to be the bottleneck & we find a very good agreement with Bell's prediction for the growth of  $E_{max}(t)$

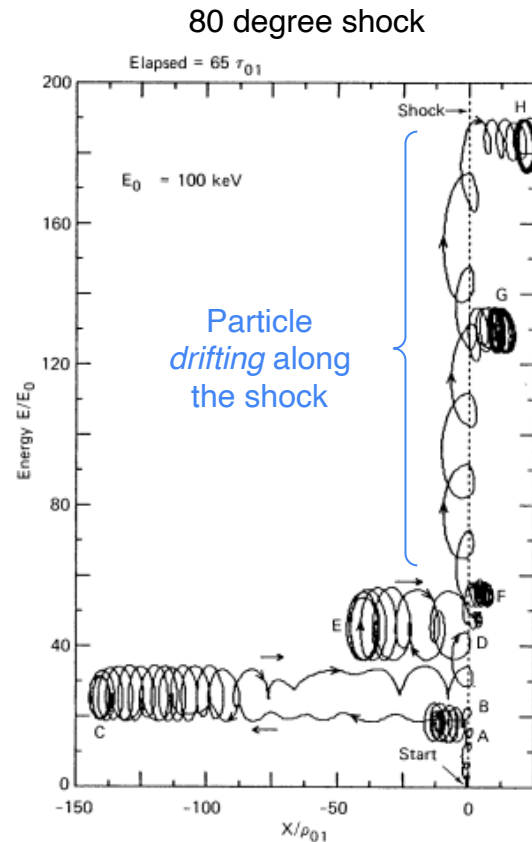
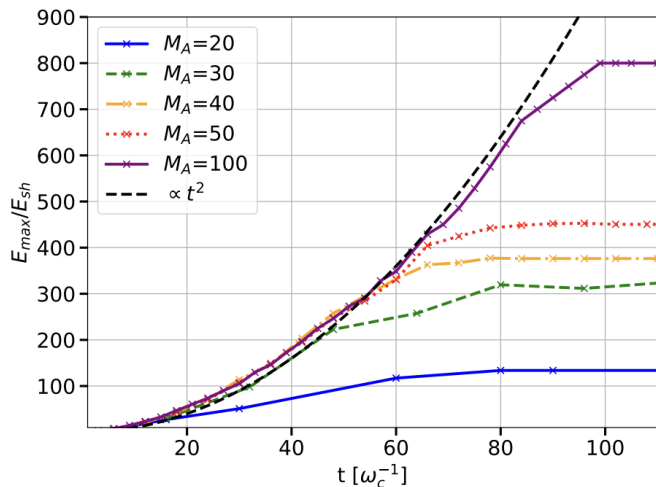


# Oblique shocks for PeV particles?

Quasi-perpendicular shocks are fast accelerators (*Kamijima et al. 2020*, *Jokipii 1982*)

Shock Drift Acceleration (SDA):  $E \propto t^2$

- Typically short-lived before advection wins
- *Orusa & Caprioli 2023* shows that the SDA maximum energy saturates at  $E_{max} \propto M_A$



Decker 1988

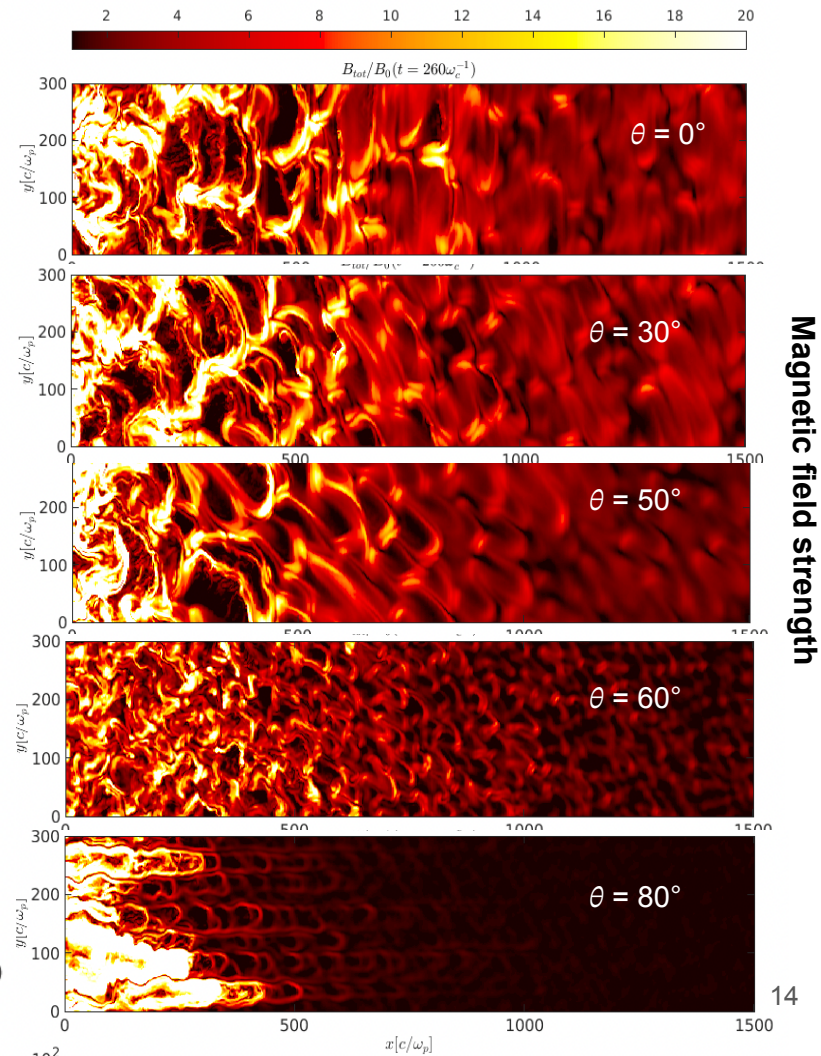
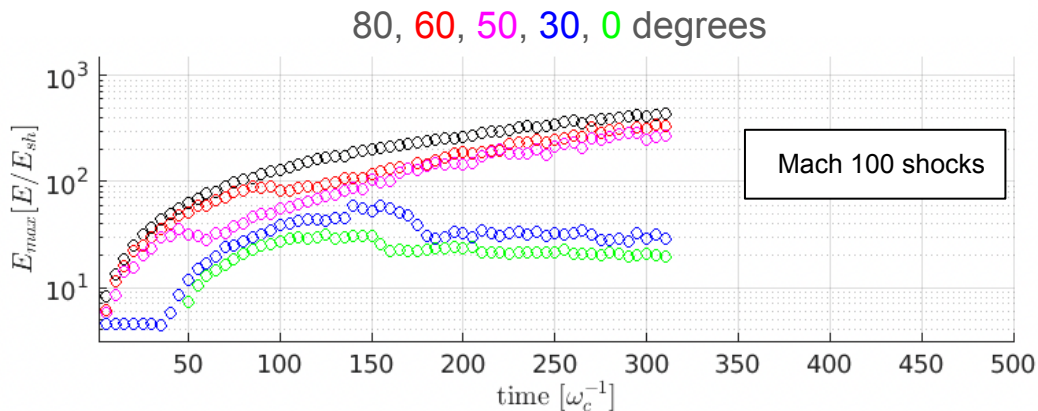
**Is it possible that Shock Drift Acceleration (SDA) will eventually transition to DSA?**

# Shock Obliquity

The size of cavities and spatial extent of turbulence is related to the number of e-foldings of the Bell Instability

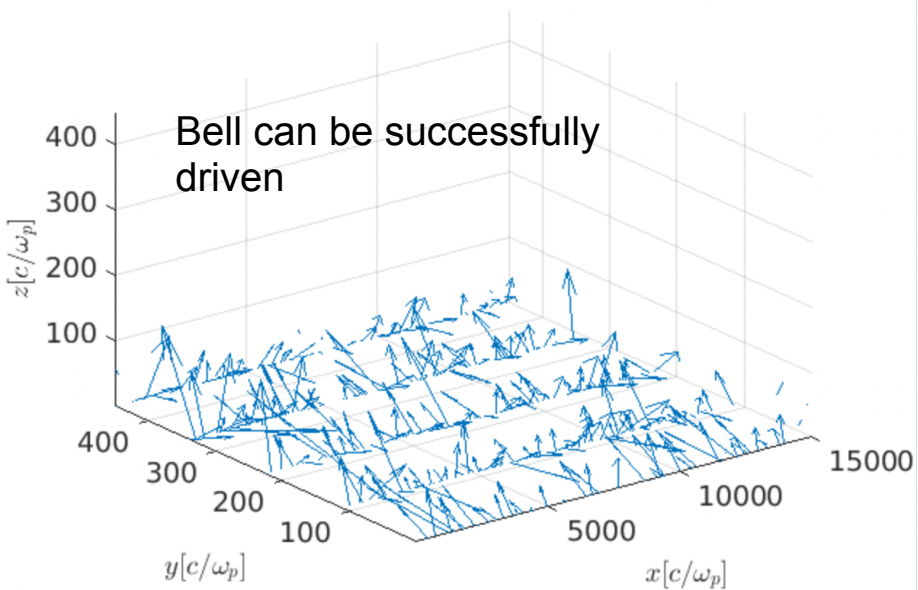
We see that **high obliquity shocks are very inefficient at triggering Bell**

However, they can still rely on Shock Drift Acceleration for brief periods of particle acceleration

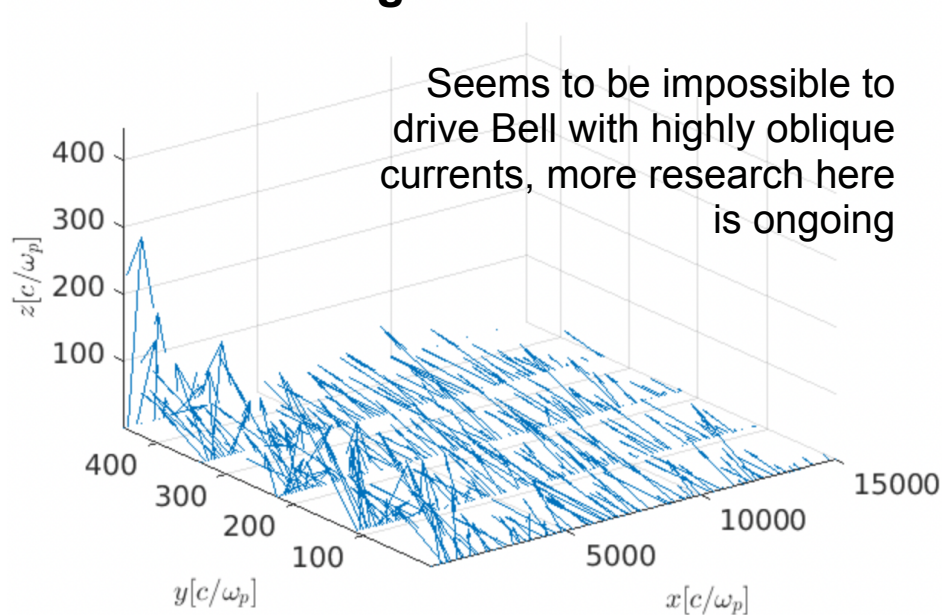


# Direction of Current Depends on Shock Obliquity

## Parallel simulation



## 50 degree simulation



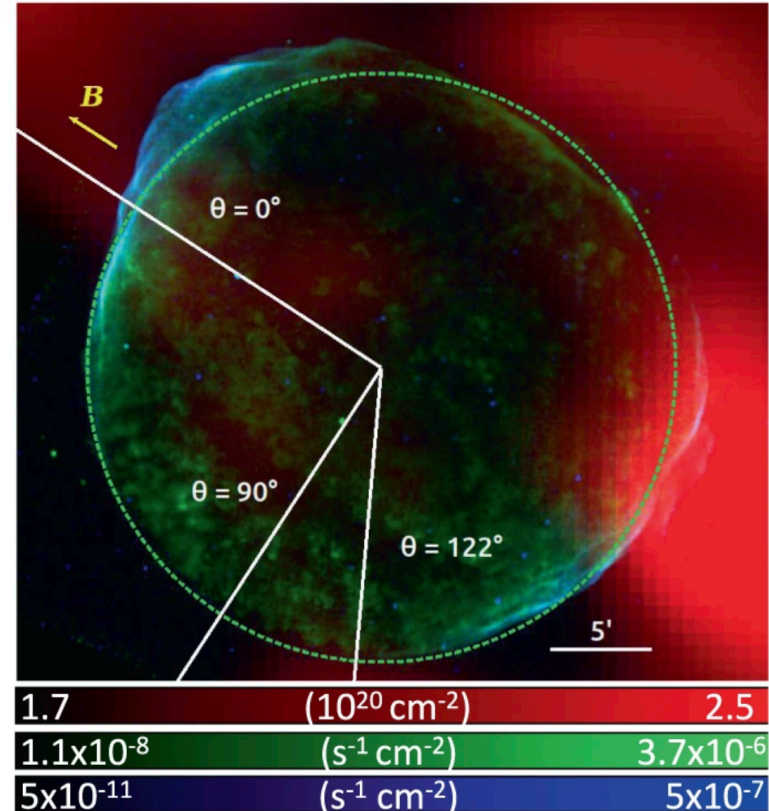
**Unlikely** that we can transition to DSA in quasi-perpendicular shocks

# Can Perpendicular Shocks be PeVatrons?

It seems **not to be possible**, which is in good agreement with observations of perpendicular shocks

In reality, CRs at highly perpendicular regions may drift around the circumference of the shock via SDA into more quasi-parallel regions

- More work remains to be done on this
- Back of the envelope: a CR could travel 10 degrees along the circumference of a remnant like SN1006 in less than 10 years
- Empirically from simulations, the maximum energy scales as  $1/\cos(\theta)$ , which is still consistent with observations



Chandra flux images of SN 1006,  
*Giuffrida et al. 2022*



# Summary

Simulations of parallel shocks seem to support the notion that the **maximum energy CRs from SNRs can only reach ~100s of TeV in energy**

To get higher energy CRs, you need strongly amplified B fields

To get strongly amplified B fields, you need the Bell Instability

To get the Bell Instability, you need to **wait**

Quasi-perpendicular shocks are still being explored, but they are unlikely to be PeVatrons

The eventual fate of CRs pushed out of perpendicular regions via SDA is still to be understood

It seems likely that we will have to rely on **other astrophysical sources** within the galaxy to provide PeV CRs like stellar clusters (*Morlino et al. 2021* , *Vieu & Reville 2022*, *Aharonian et al. 2019*).

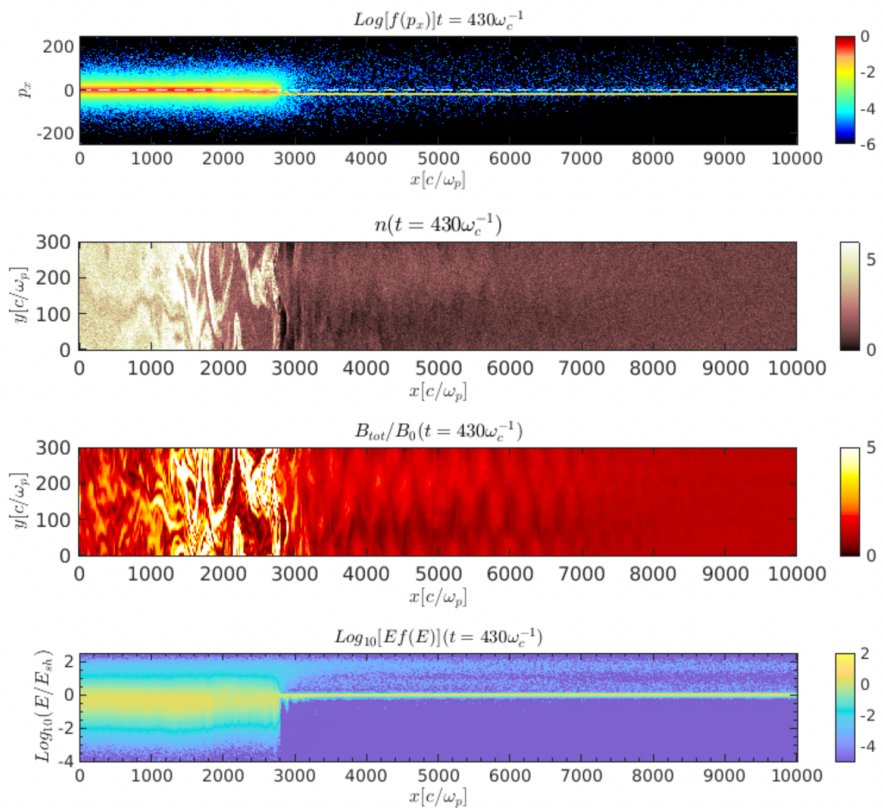
Thanks!

ευχαριστώ!

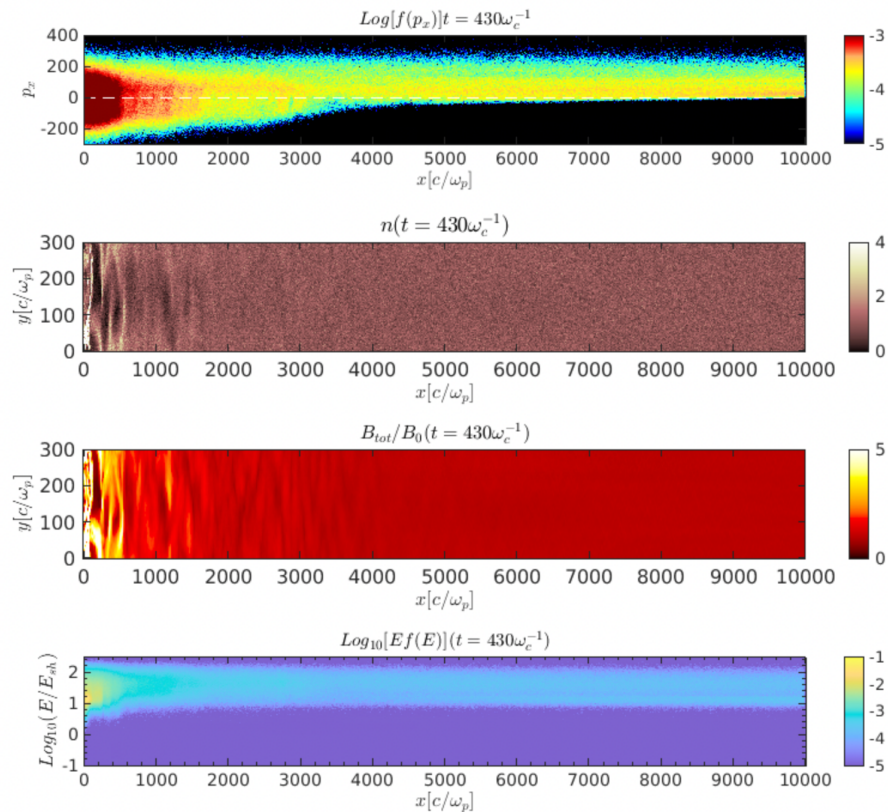
Backup

## Comparison between faux shock and real shocks

2D benchmark simulation, all particles, Ma = 20

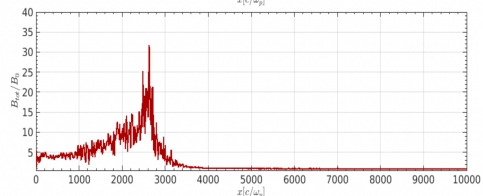
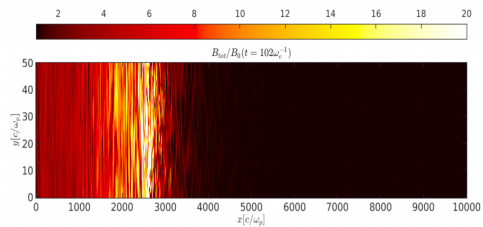


Corresponding faux shock sim, just CRs, Ma=20



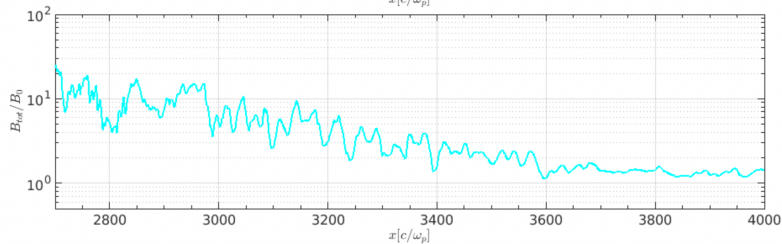
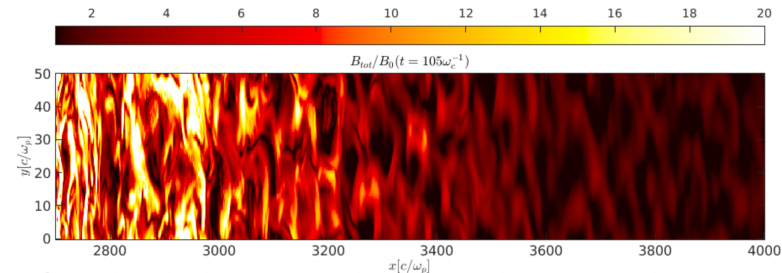
# Comparison between faux shock and 3D real shocks

80 degree, Ma = 100

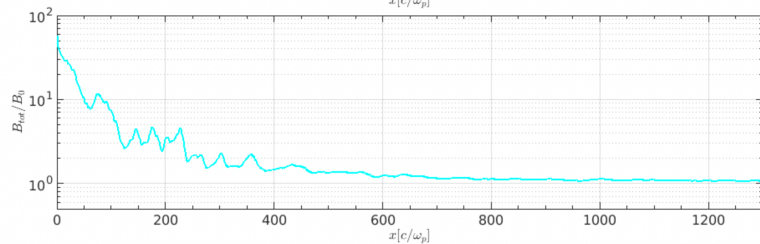
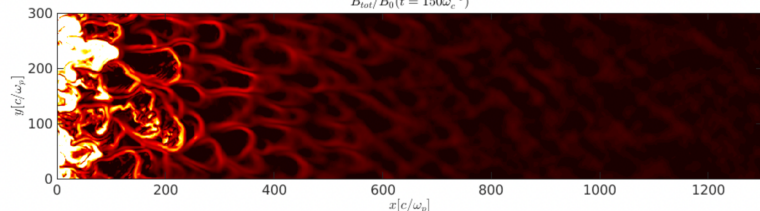
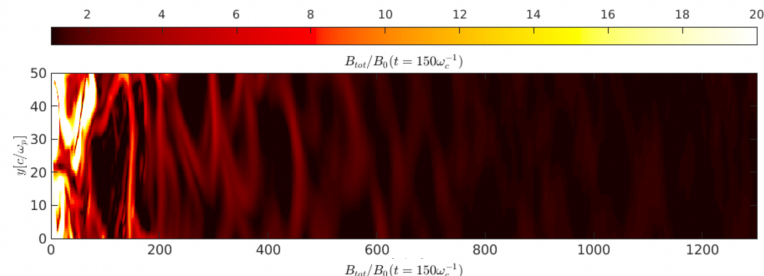


3D 80 degree sim with real shock, final timestep

zoom in ahead of shock:



Faux shock zoom in (in x and in y):



# Diffusion coefficient as a function of energy

