

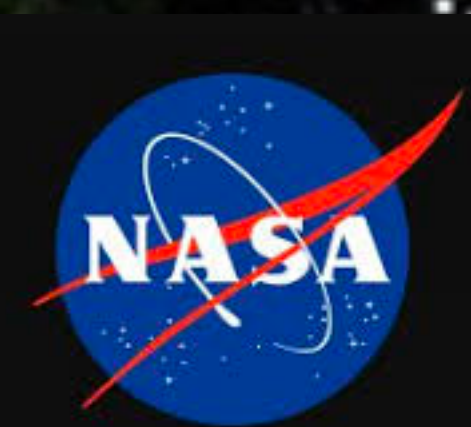
Platanias, Crete
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Particle Acceleration at SNR Shocks: Bridging Simulations and Observations

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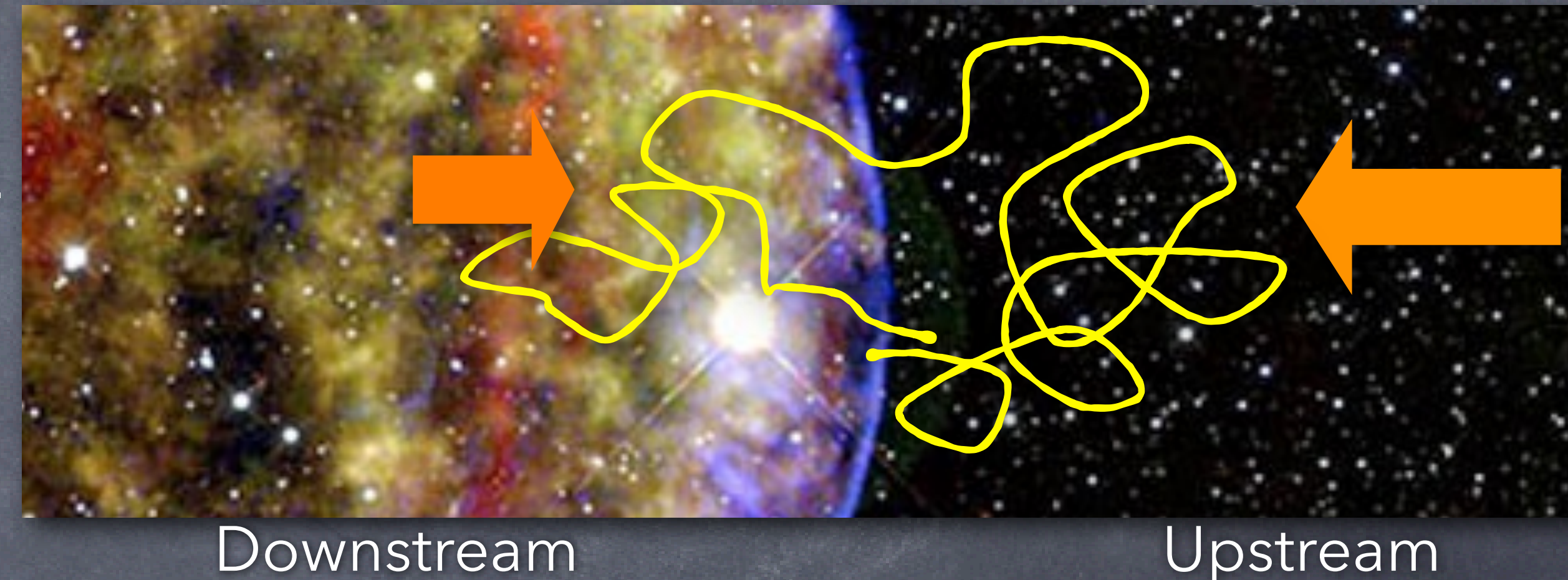


The SNR paradigm for the origin of CRs

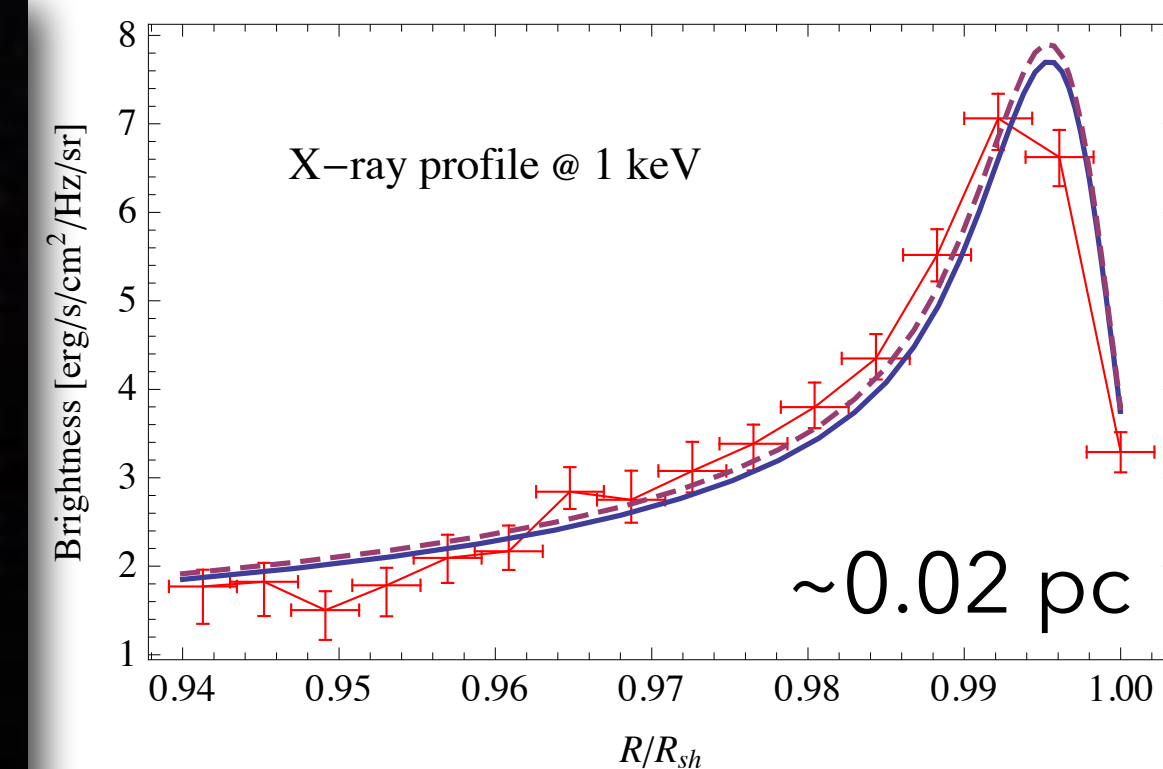
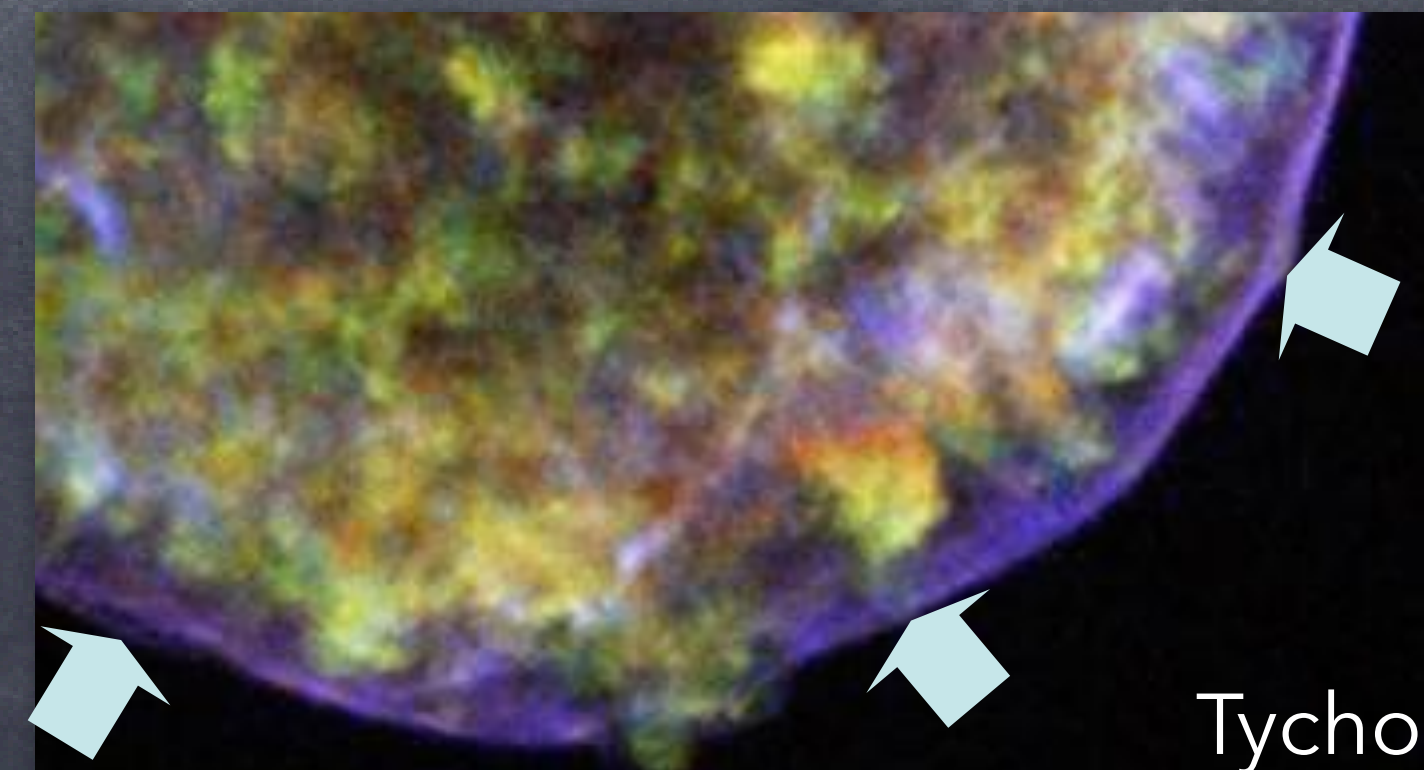
- Energetics:** $\sim 10\%$ of SN kinetic energy can account for Galactic CRs (Baade-Zwicky34)



- Mechanism:** Fermi acceleration at SNR shocks is *first-order* and produces power-laws. **Diffusive Shock Acceleration (DSA)** (Krimskii77, Axford+78, Bell78, Blandford-Ostriker78)



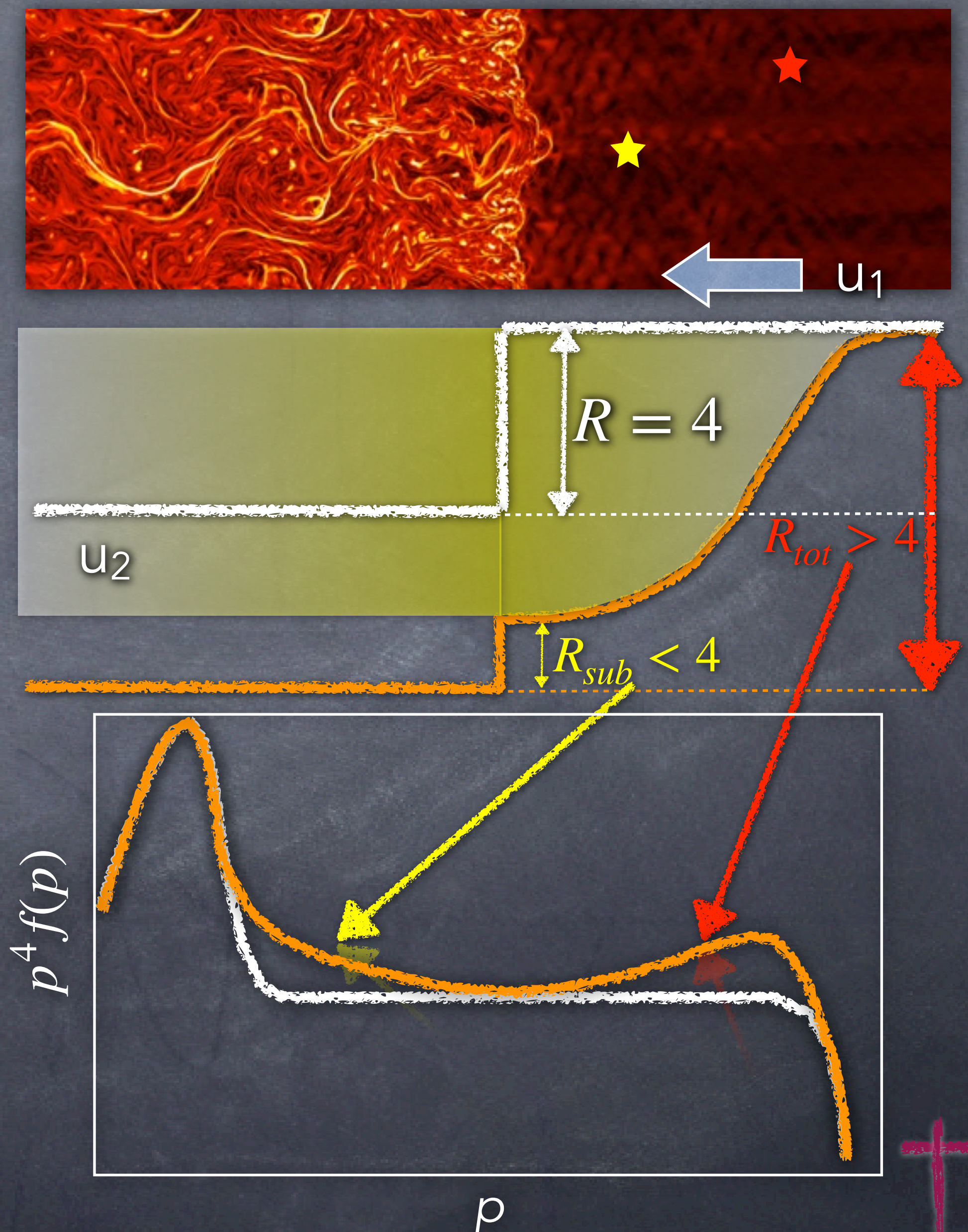
- Evidence of **B field amplification:** self-generated scattering enhances the energization rate (e.g., Bamba+05, Völk+05, Parizot+06, Morlino+12, Ressler+14, etc)



Non-Linear Diffusive Shock Acceleration

- DSA yields *momentum* power laws $f(p) \propto 4\pi p^2 p^{-q}$
- The **slope** q depends only on the shock **compression**

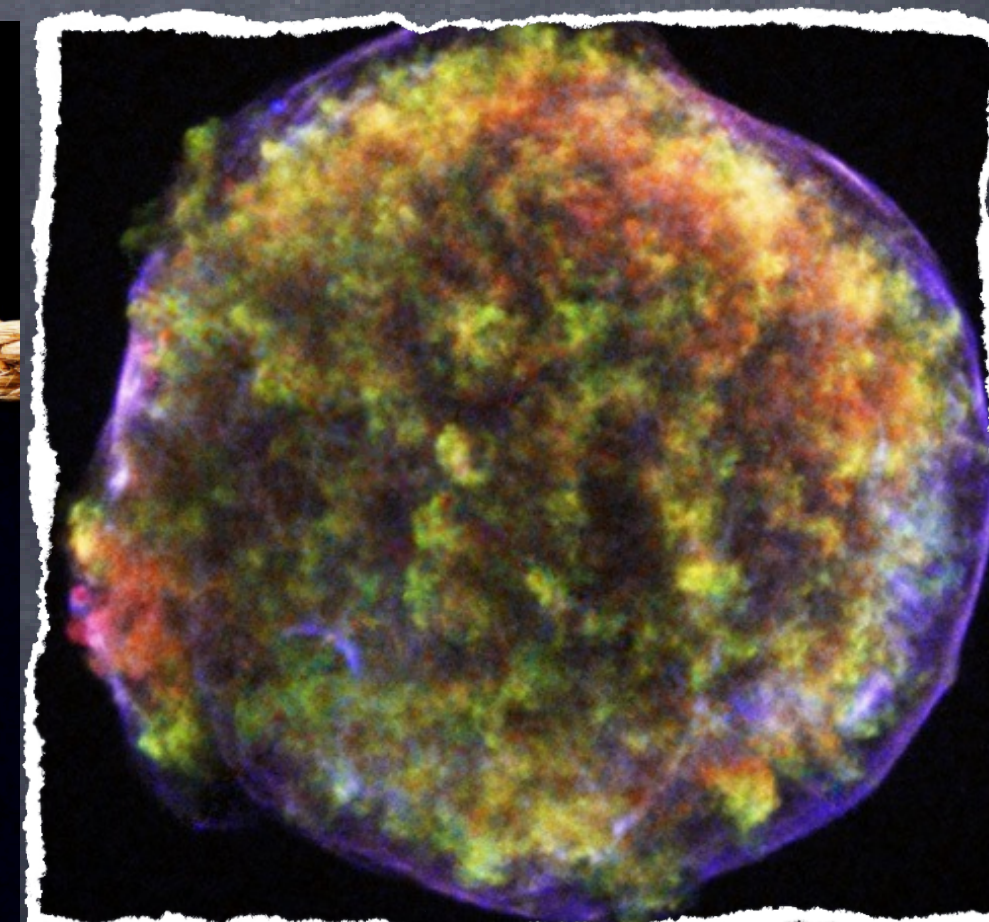
$$q = \frac{3R}{R-1}; \quad R = \frac{\gamma+1}{\gamma-1} \simeq 4; \rightarrow q = 4 \text{ for strong shocks}$$
- The CR pressure makes the **adiabatic index** γ smaller and induces a shock **precursor**
- Particles "feel" different compression ratios: spectra should become **concave**
- If **acceleration is efficient**, high-energy particles feel $R_{tot} > 4$ and their spectra must be flat, i.e., $q < 4$



(e.g., Jones-Ellison91, Malkov-Drury01 for reviews)

Theory vs Observations

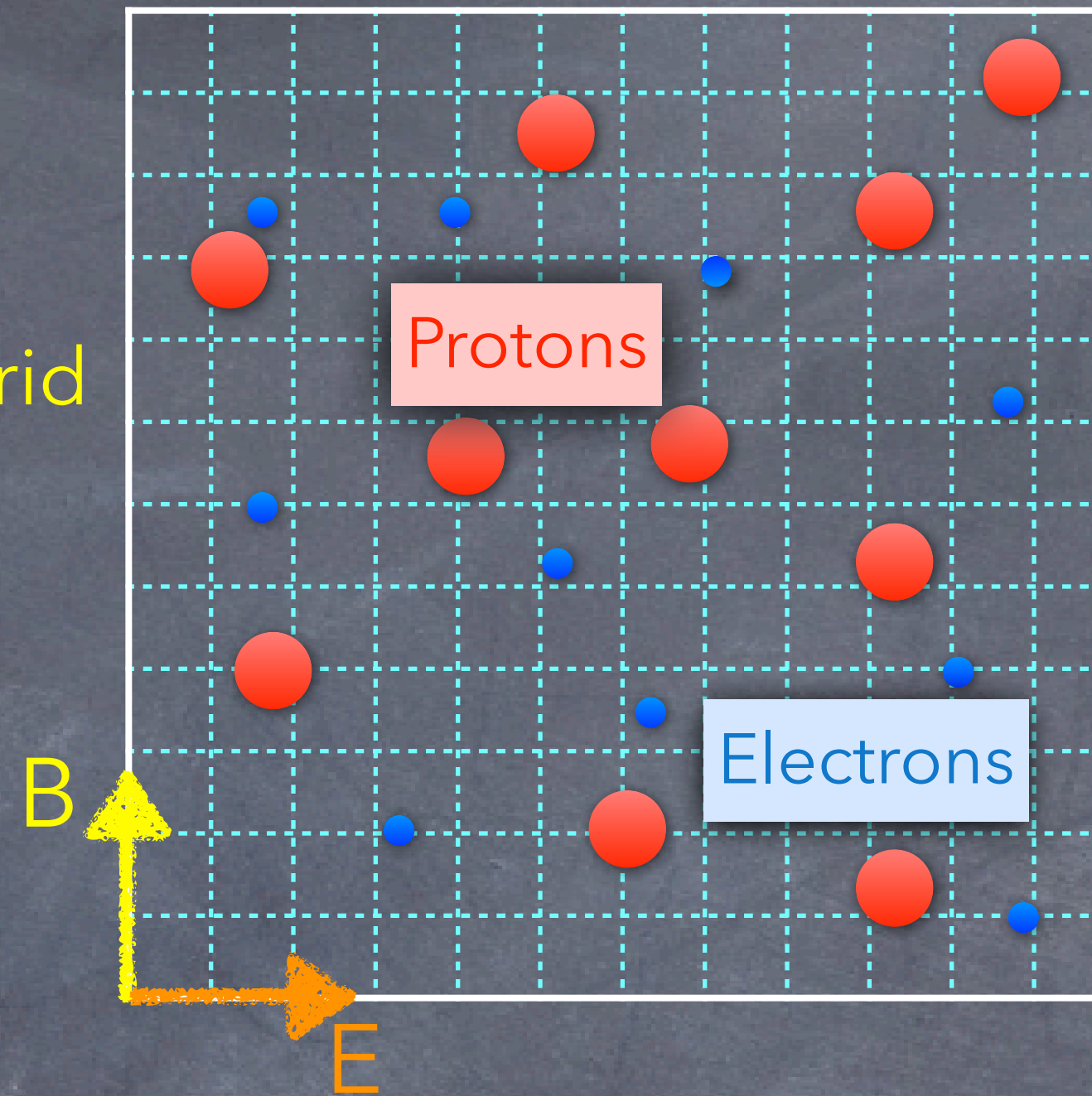
- **Efficient DSA** should return:
 - Compression ratios $R > 4$;
 - CR spectra flatter than p^{-4} (flatter than E^{-2} for relativistic particles)
- **Observations**, instead, point to significantly steeper spectra:
 - Hadronic γ -rays from historical and middle-age SNRs: $q \sim 4.3 - 4.7$ (e.g., Caprioli11,12; Aharonian+19);
 - Synchrotron emission from radio SNe: $q \sim 5$ (e.g., Chevalier-Fransson06, Bell+11);
 - Propagation of Galactic CRs suggests source spectra with $q \sim 4.3 - 4.4$ (e.g., Blasi-Amato11a,b; Evoli+19).



Astroplasmas from first principles

Full-PIC approach

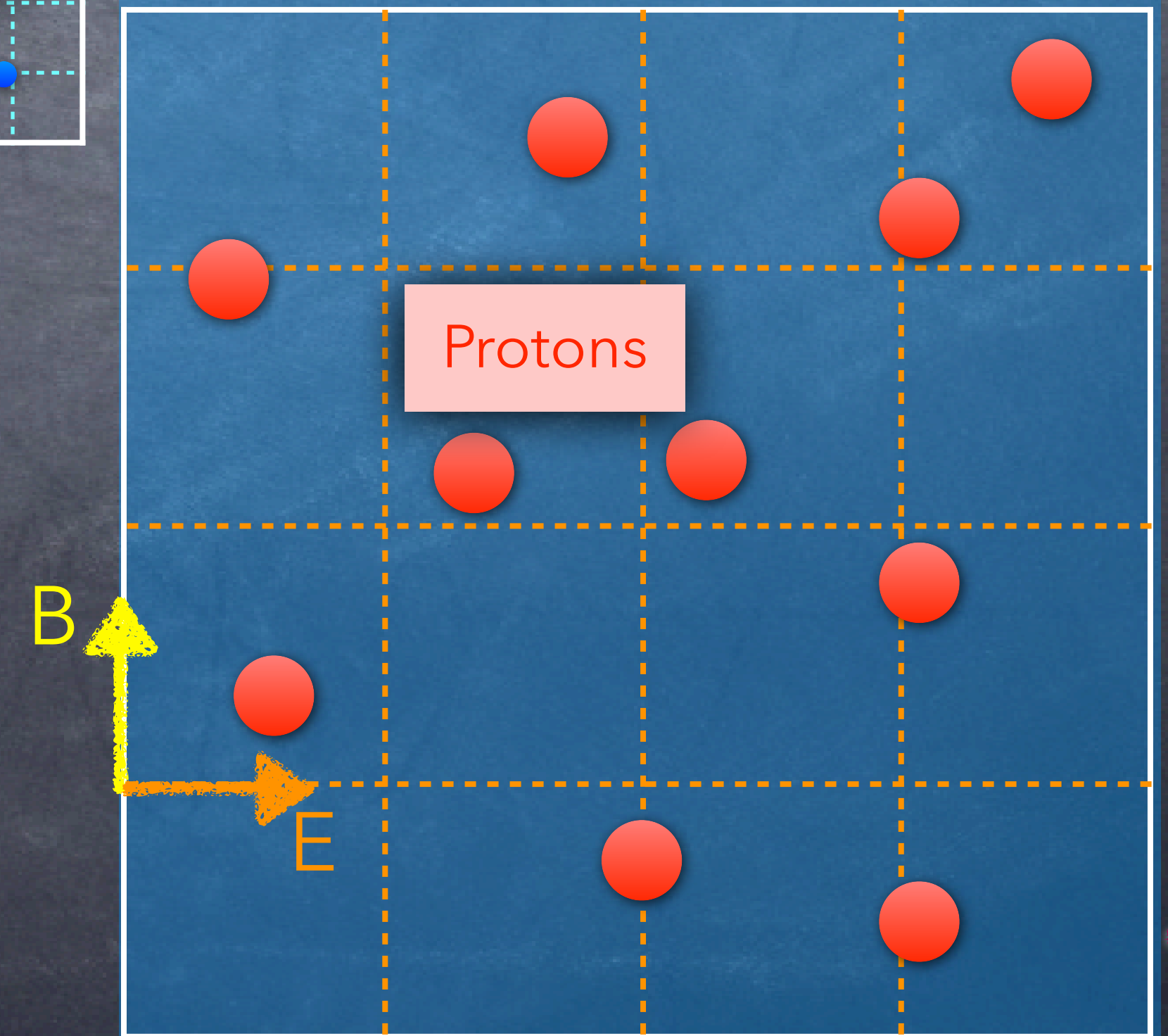
- Define electromagnetic fields on a **grid**
- Move particles via **Lorentz force**
- Evolve fields via **Maxwell equations**
- Computationally very challenging!



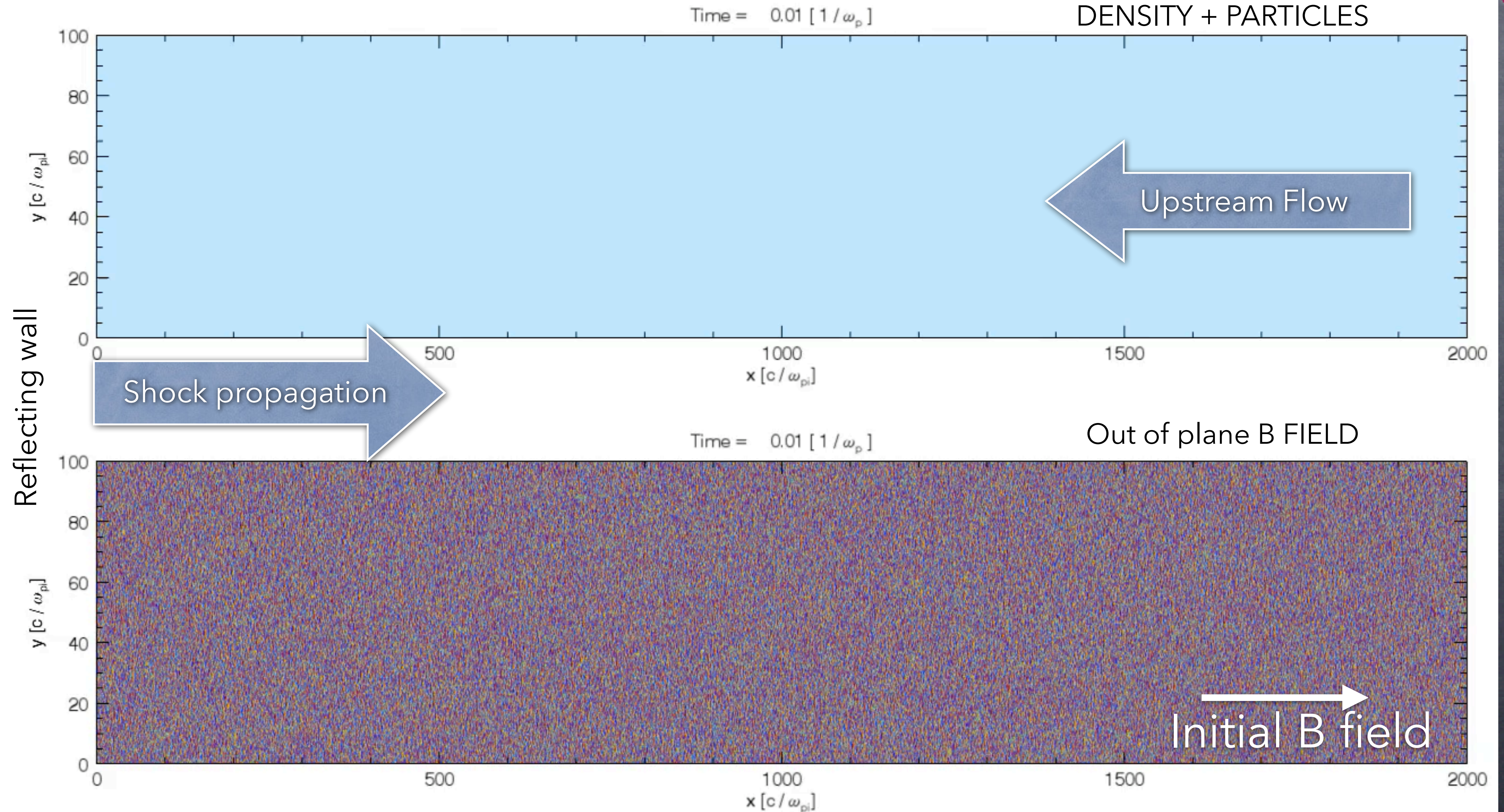
Hybrid approach: Fluid **electrons** - Kinetic **protons**

(Winske & Omid; Burgess et al., Lipatov 2002; Giacalone et al. 1993, 1997, 2004-2013; DC & Spitkovsky 2013-2015, Haggerty & DC 2019...)

- massless electrons for more **macroscopical** time/length scales

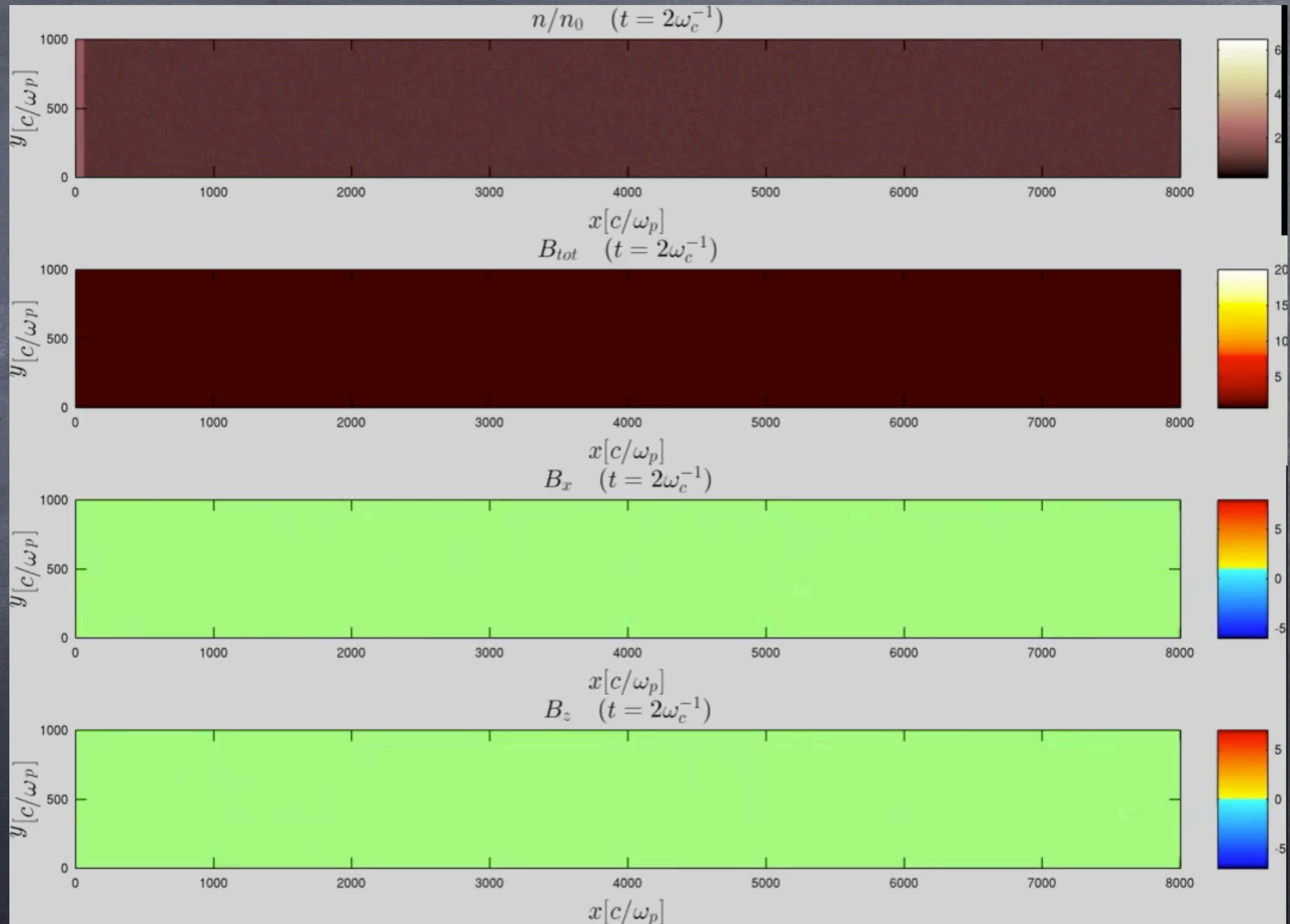


Hybrid Simulations of Collisionless Shocks



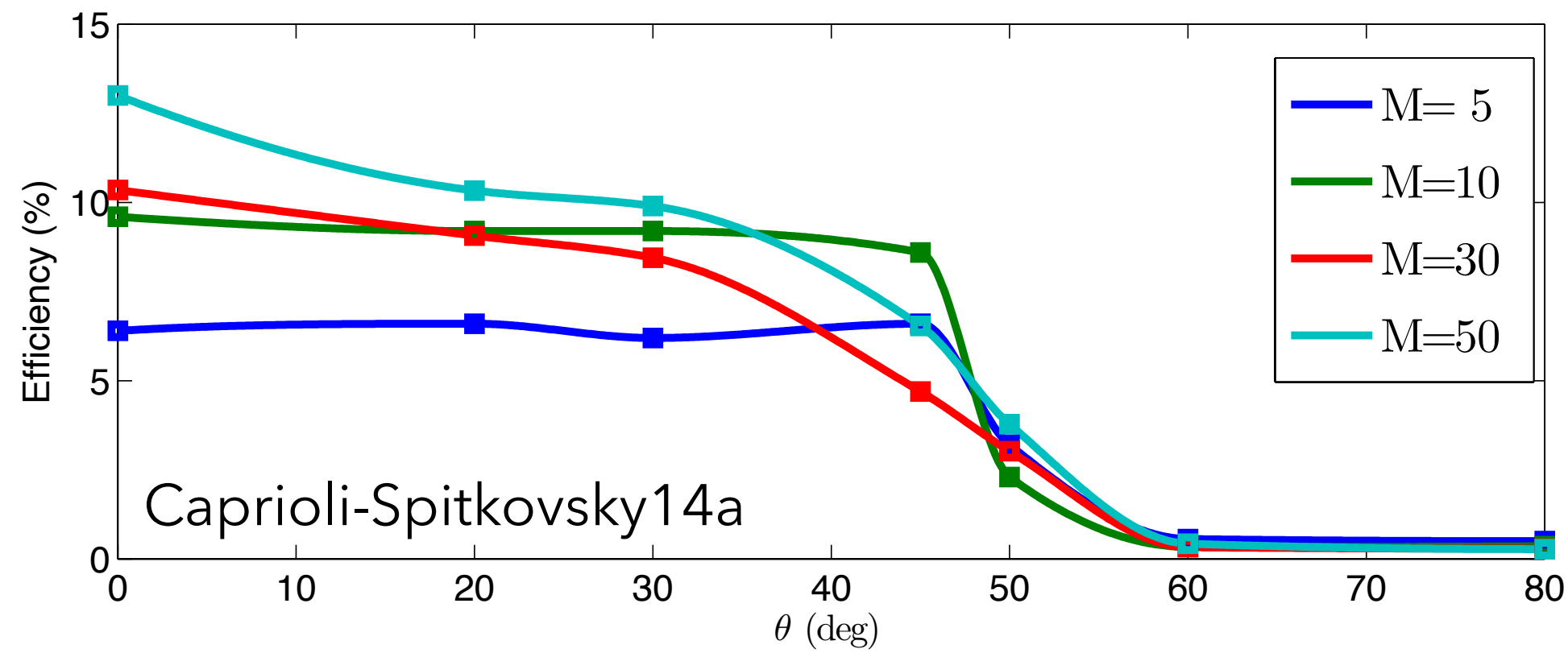
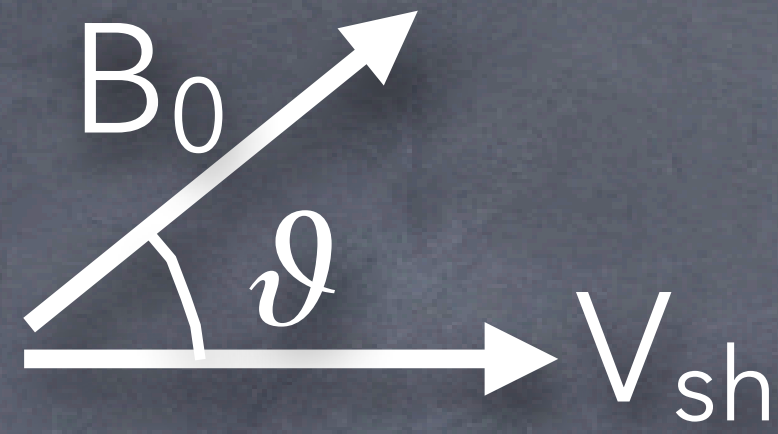
CR-driven Magnetic-Field Amplification

Initial B field
 $M_s = M_A = 30$

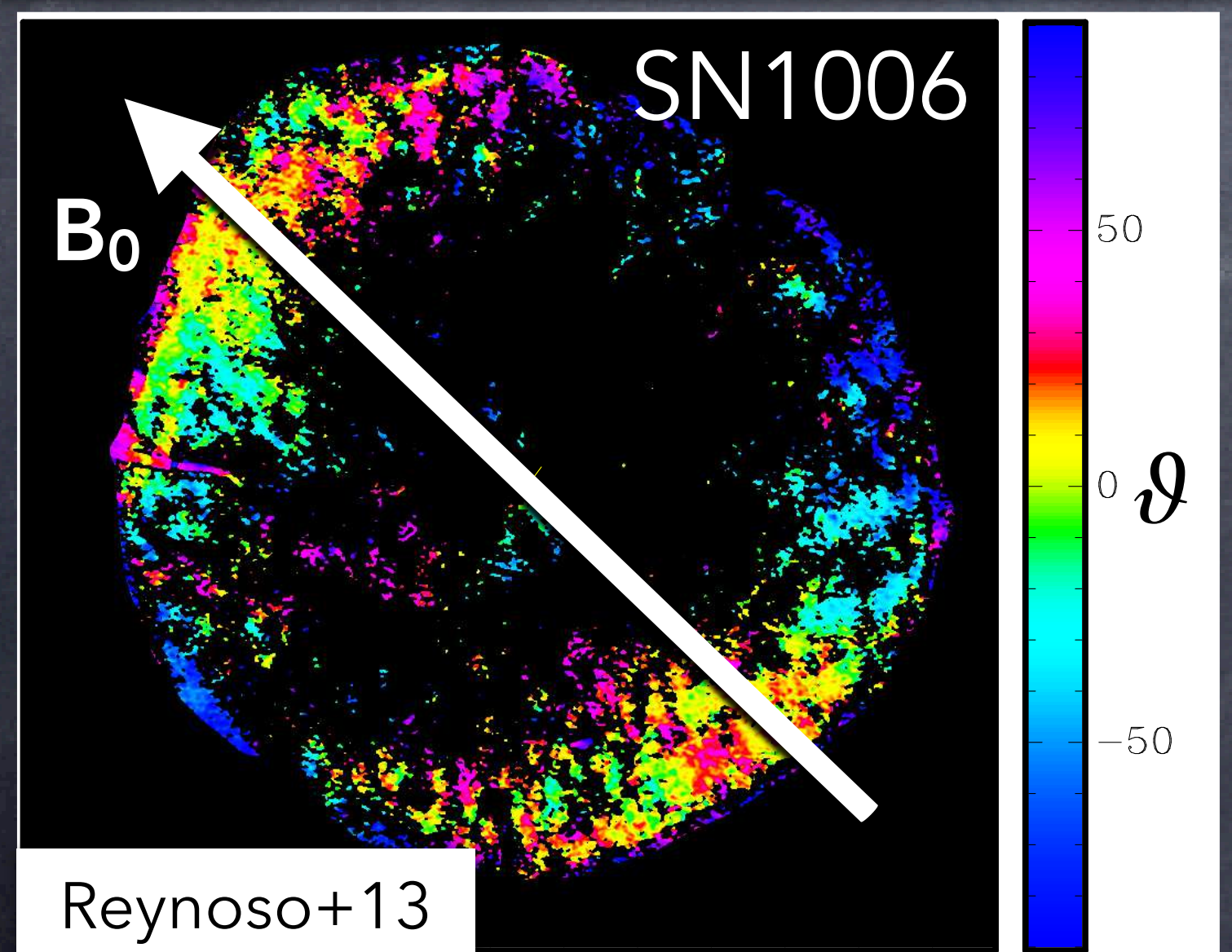
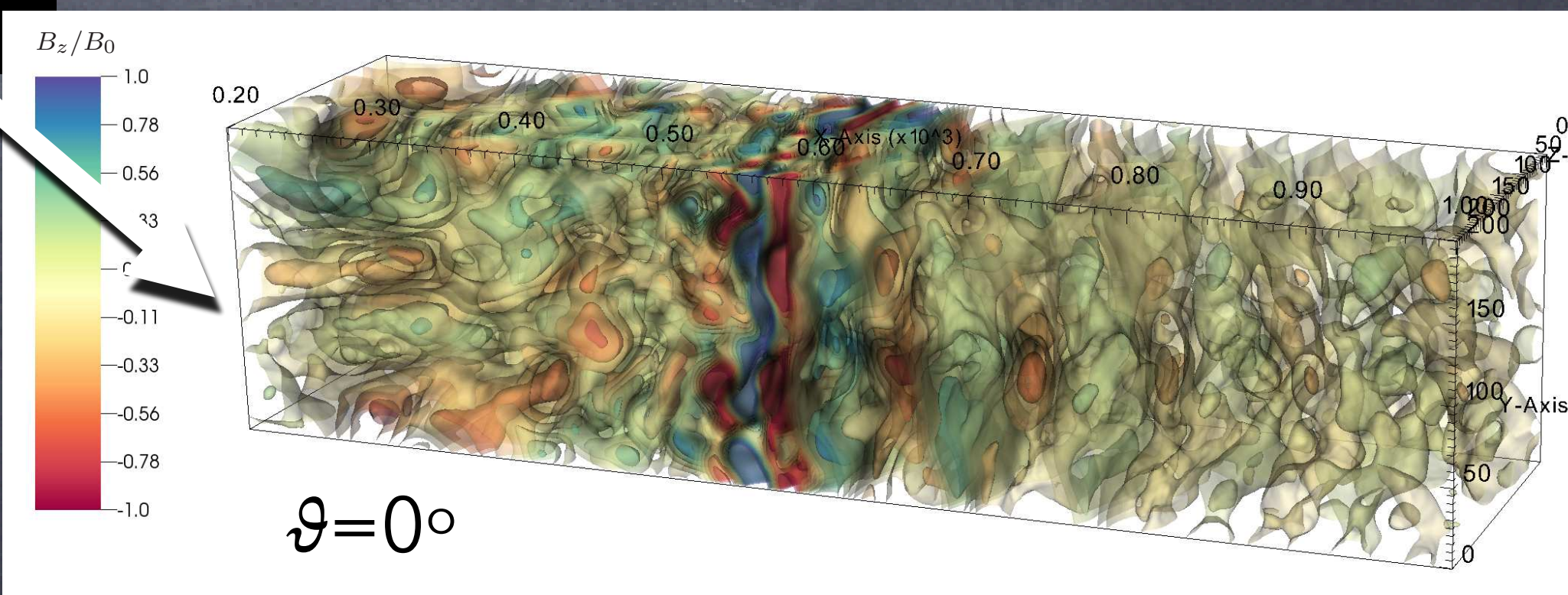
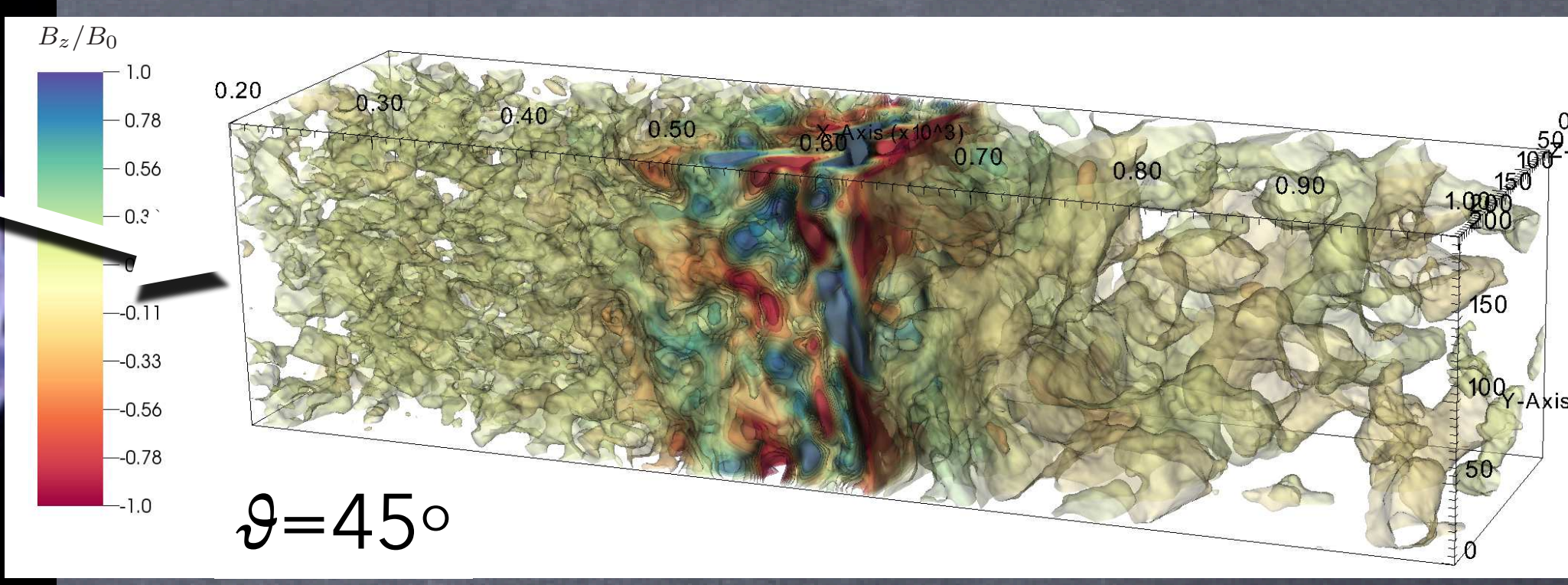
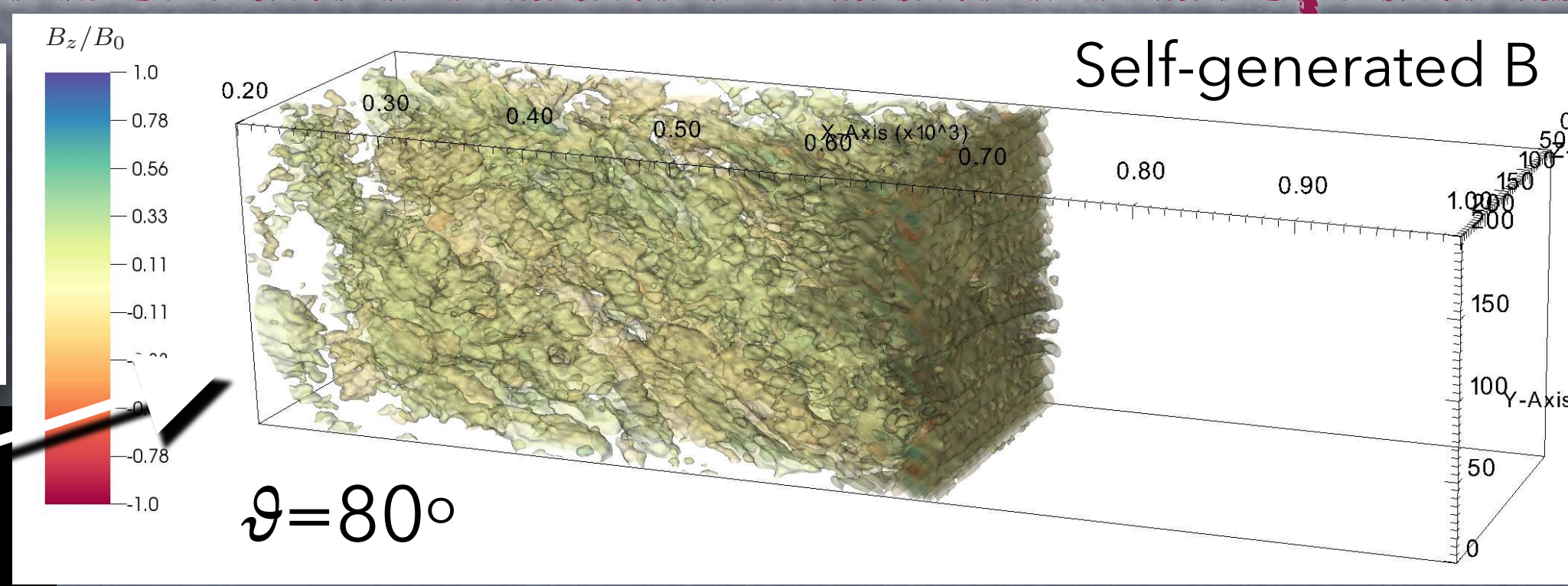
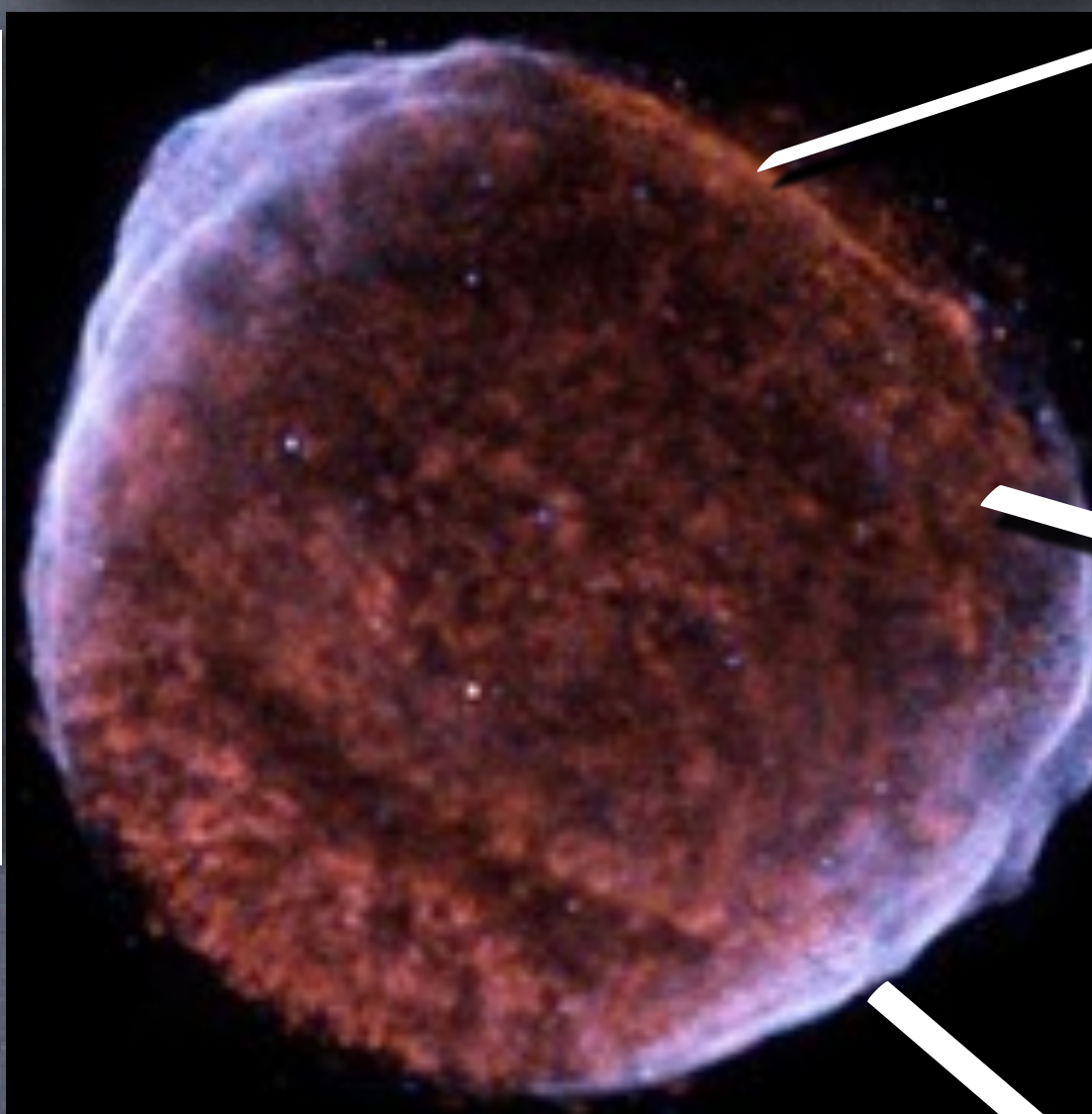


DSA Efficiency

Acceleration depends on the shock **inclination**



X-ray emission:
red=thermal
white=synchrotron



B amplification and ion acceleration where the shock is **parallel**

Caprioli-Spitkovsky14a,b,c

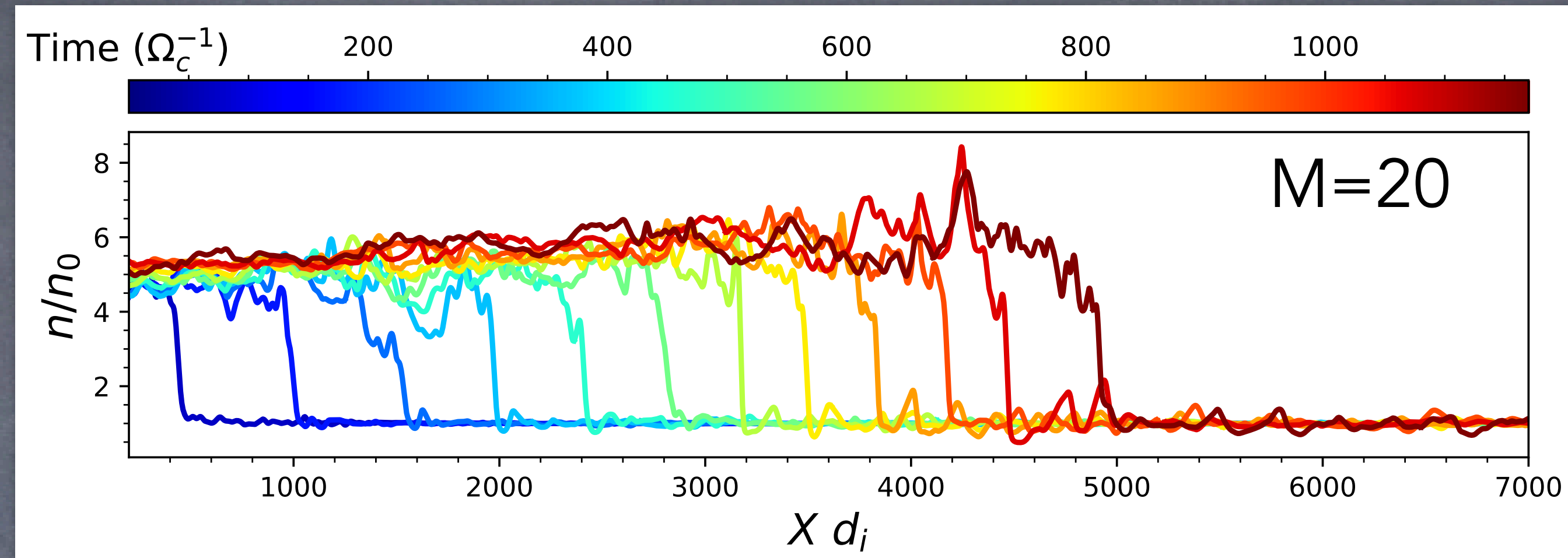
CR-modified Shocks: Enhanced compression!

- Hybrid simulations (Haggerty-Caprioli20)

- Efficiency $\lesssim 15\%$ at parallel shocks

- Formation of upstream precursor

- R increases with time, up to ~ 6



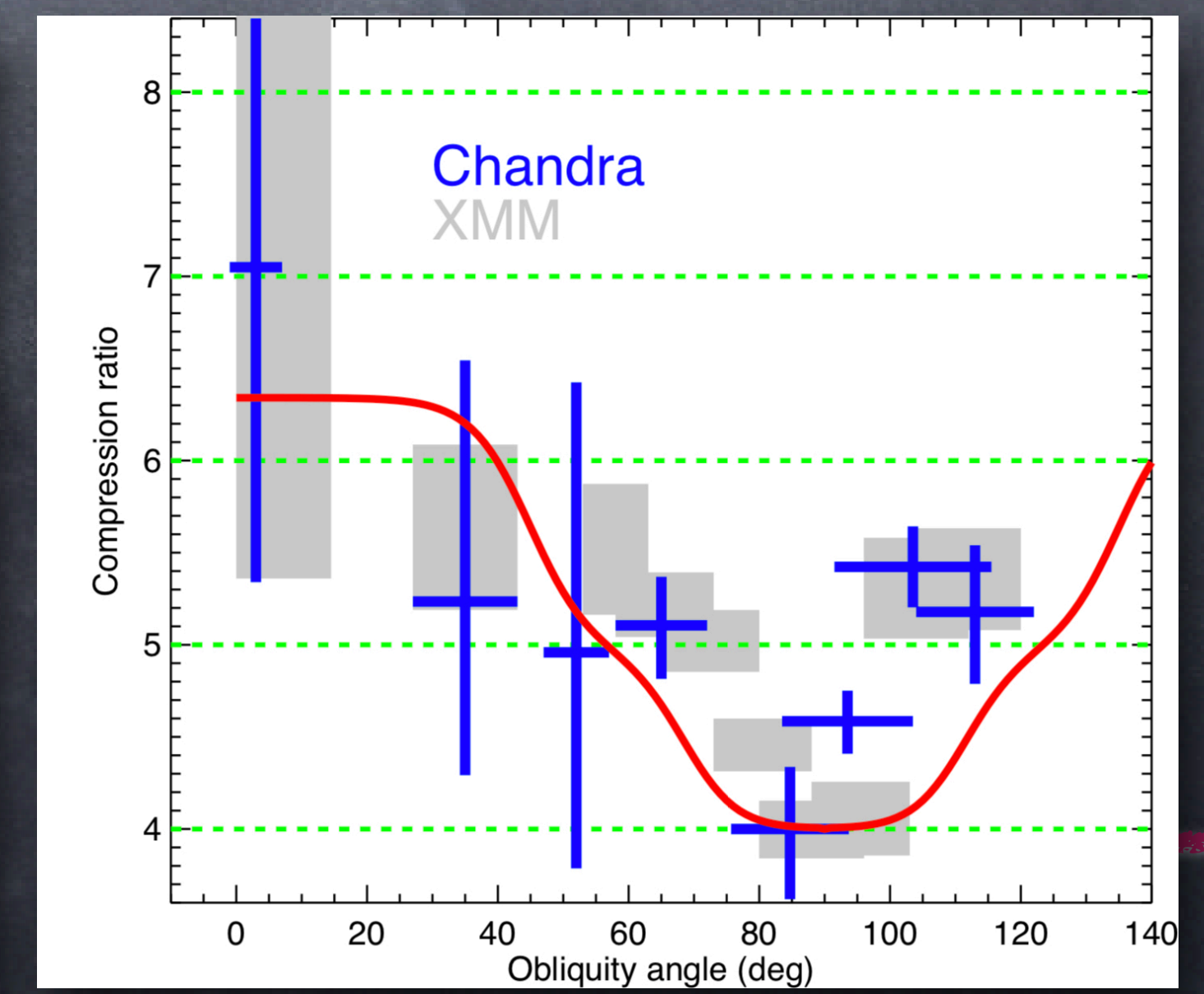
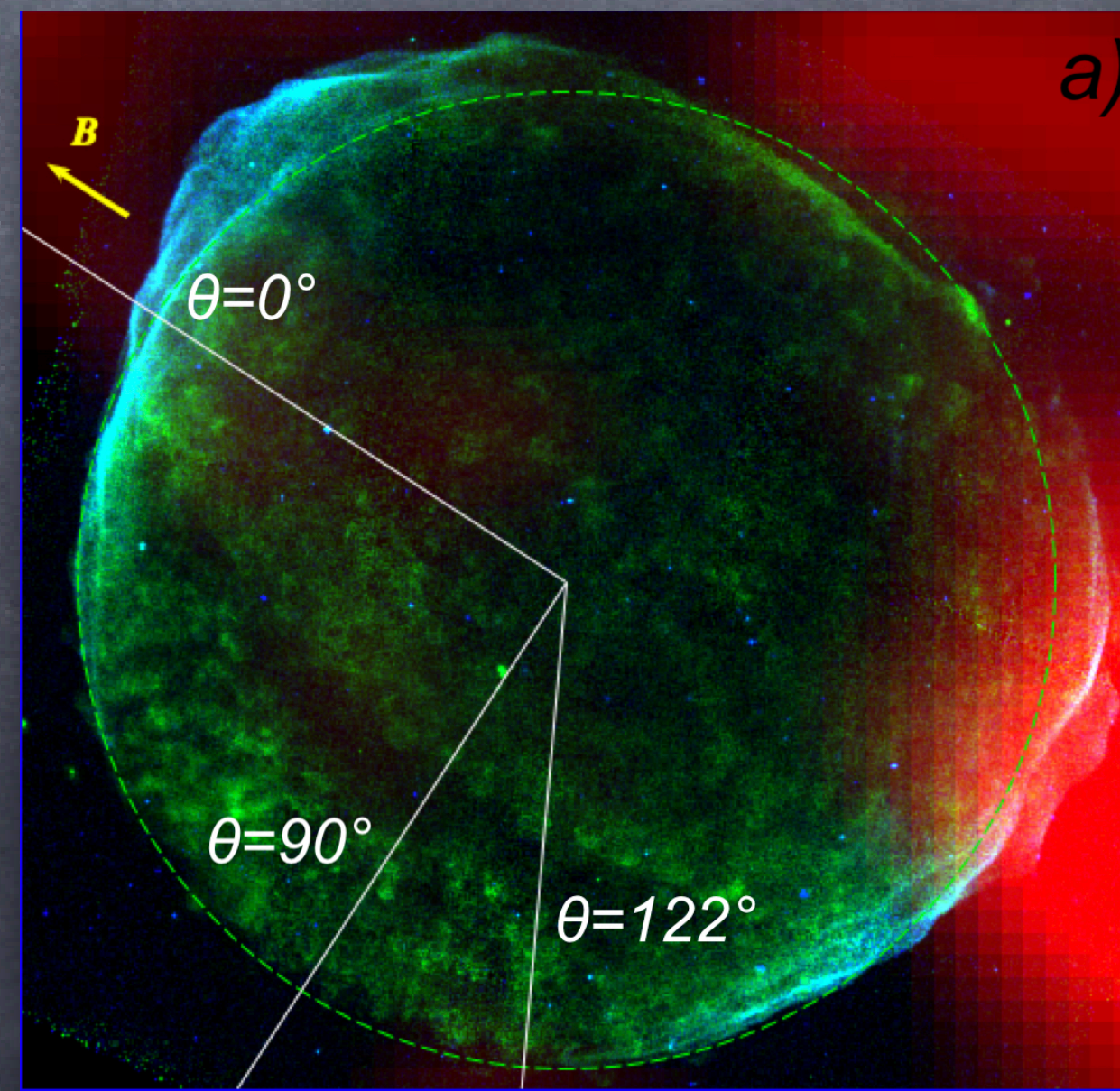
- $R \sim 6 - 7$ inferred in **Tycho** (Warren+05). In **SN1006**: $R \sim 4 - 7$, modulated with the azimuth/shock inclination (Giuffrida+21)

- If $R \simeq 7 \rightarrow q_{\text{expected}} \simeq 3.5$

- Tycho: radio to γ -ray observations:

$$q_{\text{inferred}} \simeq 4.3$$

A challenge to DSA theory!



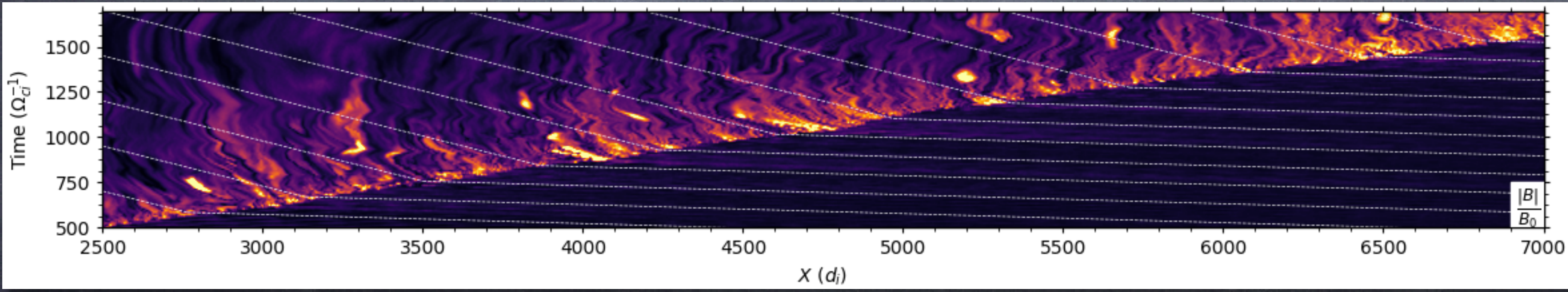
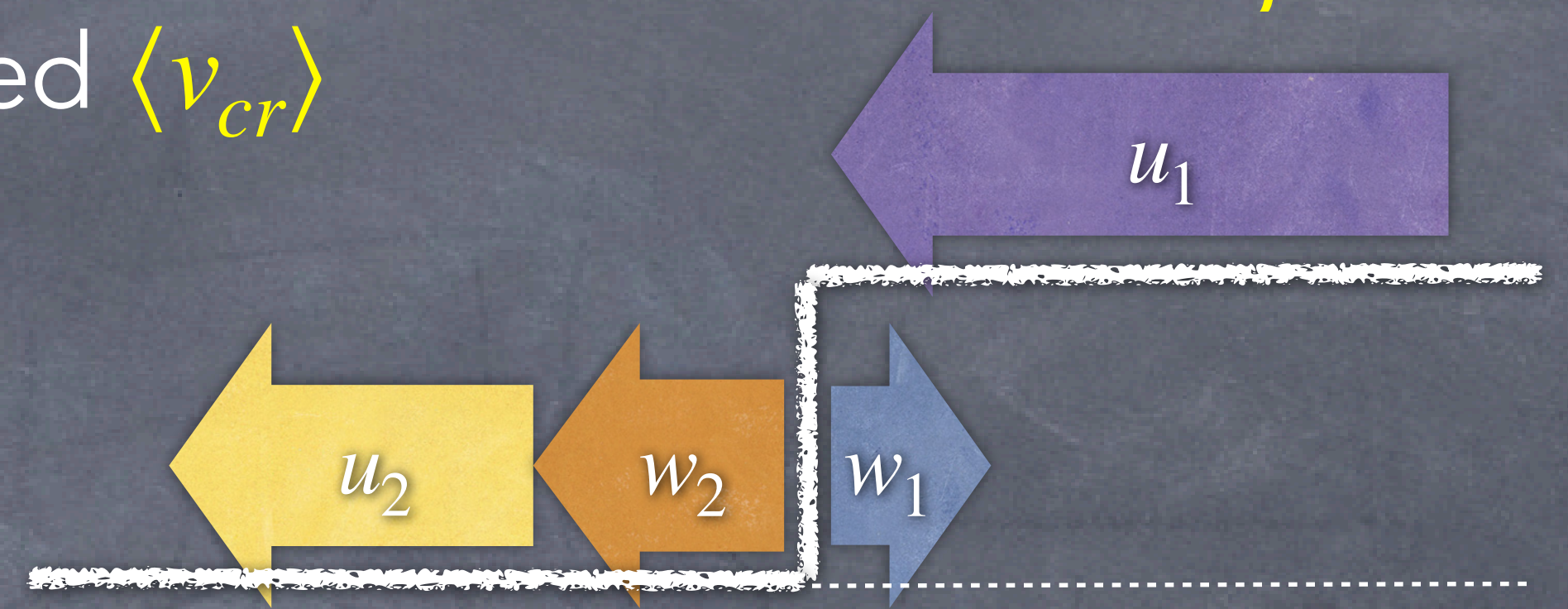
The Role of Amplified Magnetic Fields

- CRs feel an **effective** compression $R_{cr} = \frac{u_1 + w_1}{u_2 + w_2}$; $w = \text{wave speed} \approx v_A = \frac{B}{4\pi\rho}$

- We can measure both w and the effective CR speed $\langle v_{cr} \rangle$

- Upstream:** $w_1 \simeq -v_{A,1}(\delta B_1) \ll u_1$

- Downstream:** $\langle v_{cr} \rangle \simeq w_2 \simeq +v_{A,2}(\delta B_2) \equiv \alpha u_2$



Haggerty-Caprioli20

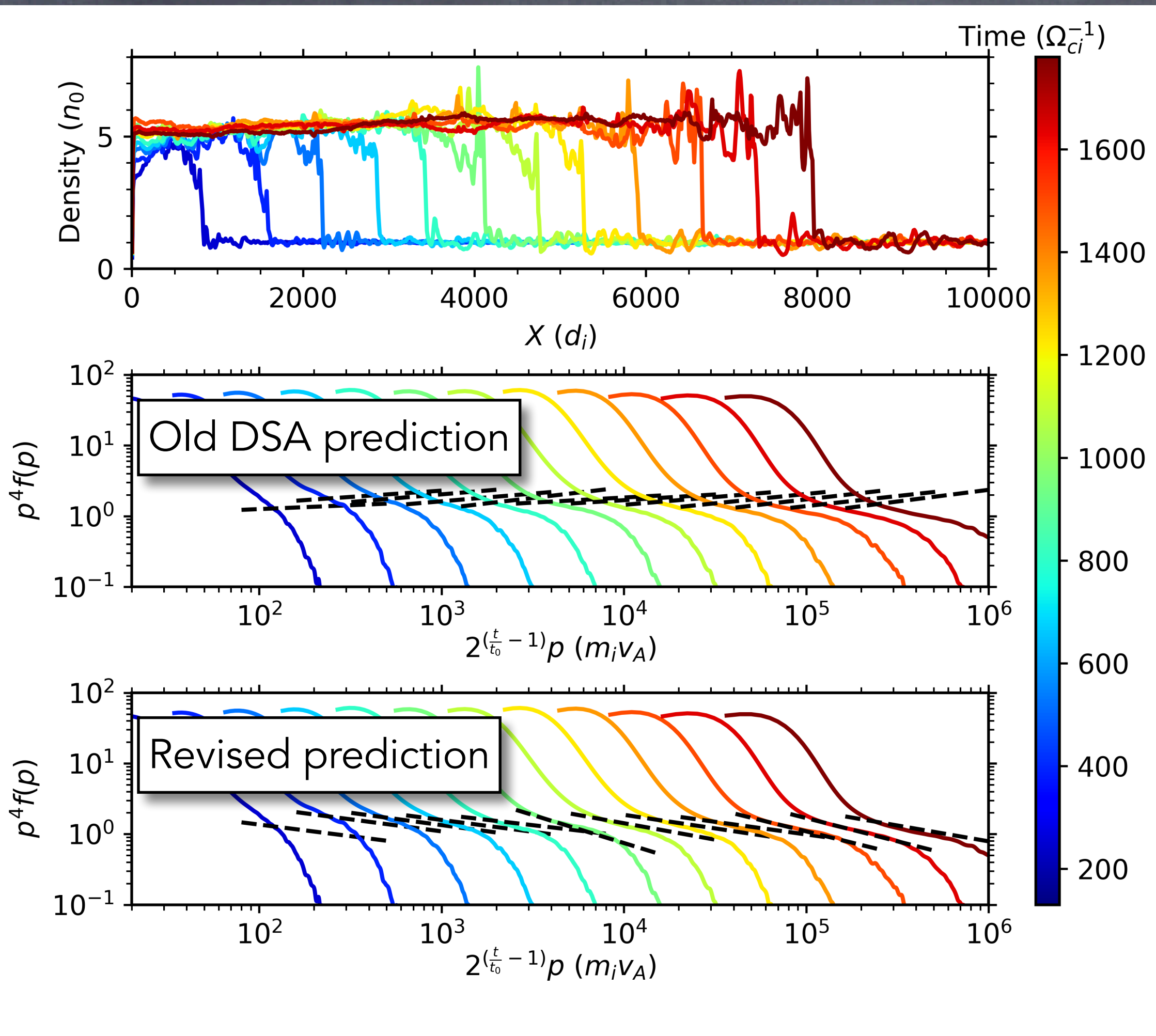
- B fields (and hence CRs) **drift** downstream with respect to the thermal gas

- First evidence of the formation of a **postcursor**

$$R_{cr} \simeq \frac{u_1}{u_2(1 + \alpha)} < R_{gas}$$

- CRs *feel* a compression ratio *smaller* than the gas

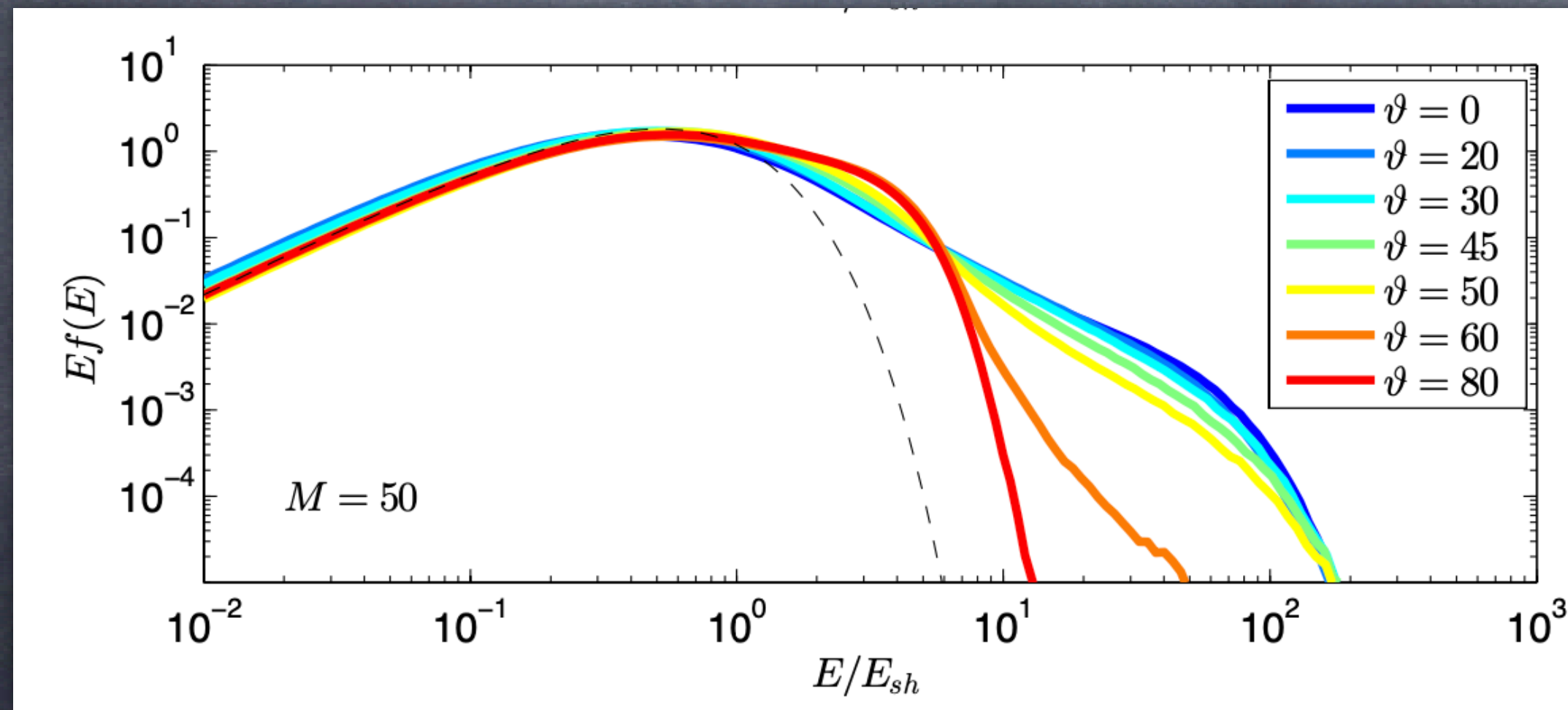
A Revised Theory of Diffusive Shock Acceleration



- With the **effective** compression felt by CRs
- $$q = \frac{3R_{cr}}{R_{cr} - 1} = \frac{3R_{gas}}{R_{gas} - 1 - \alpha} > q_{DSA}$$
- CRs feel $R_{cr} < R_{gas}$: the power-law index is *not universal, but depends on B field*
- Ab-initio* explanation for the **steep spectra observed** in SNRs, radio SNe, CRs...
 - Diesing-Caprioli21; See also R. **Diesing's** talk
- In a multi-wavelength fit B strength and particle slope are not independent!

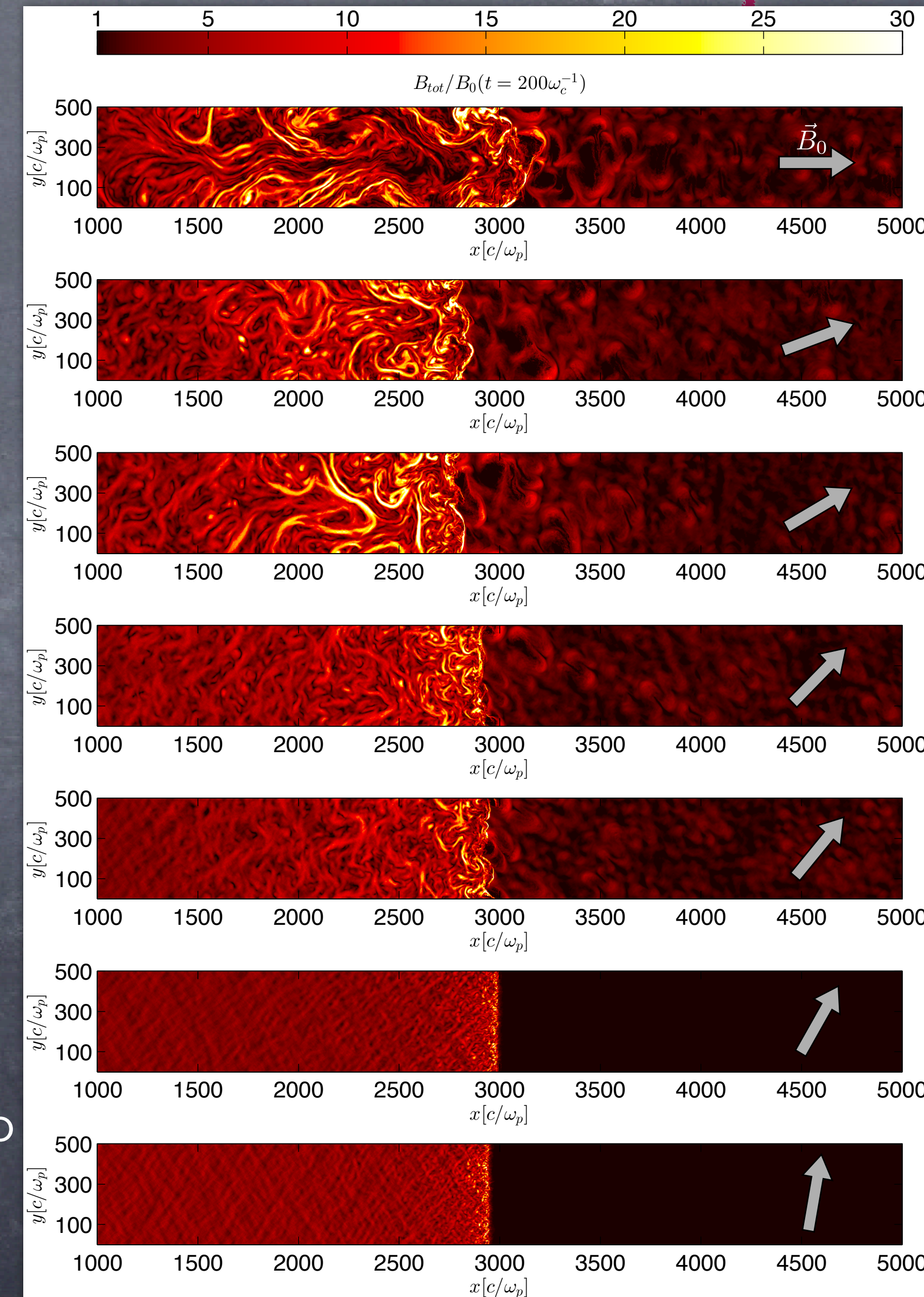
Oblique Shocks

- Oblique shocks are good accelerators but **bad ion injectors** (Jokipii82, Giacalone+00, Giacalone05, Caprioli+15)
- Is there a **critical magnetization** ($\propto 1/M_A^2$) below which \mathcal{D} becomes *irrelevant*?
- No evidence in 2D hybrid sims **w/o CR or B seeds**



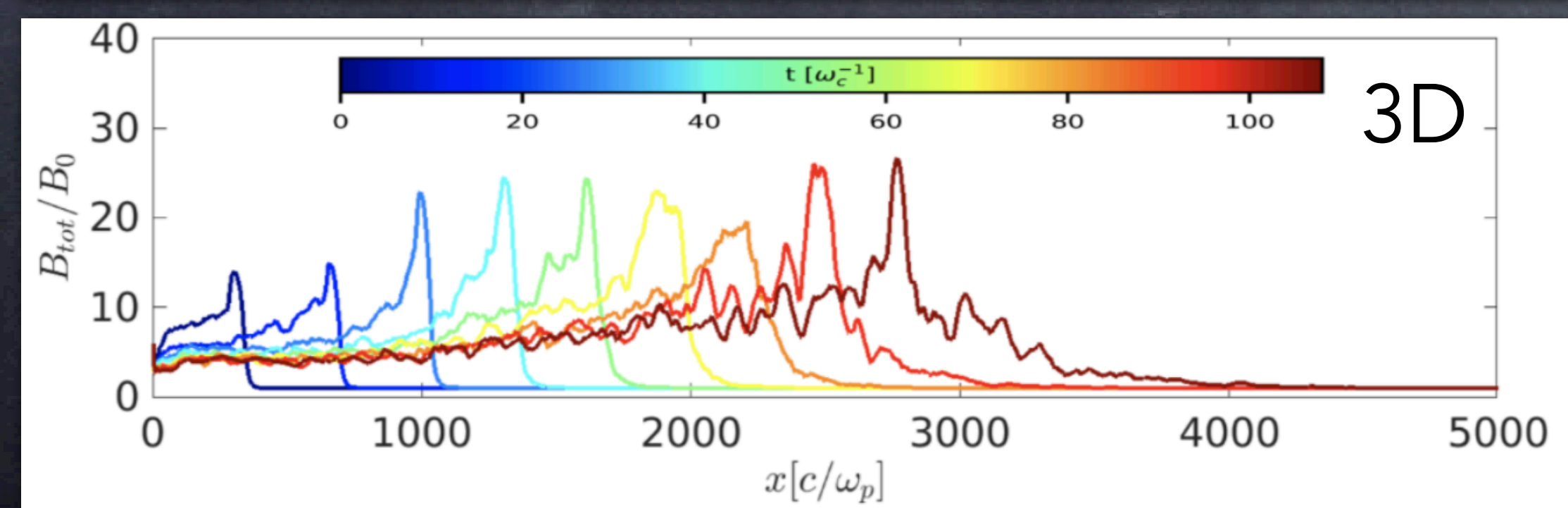
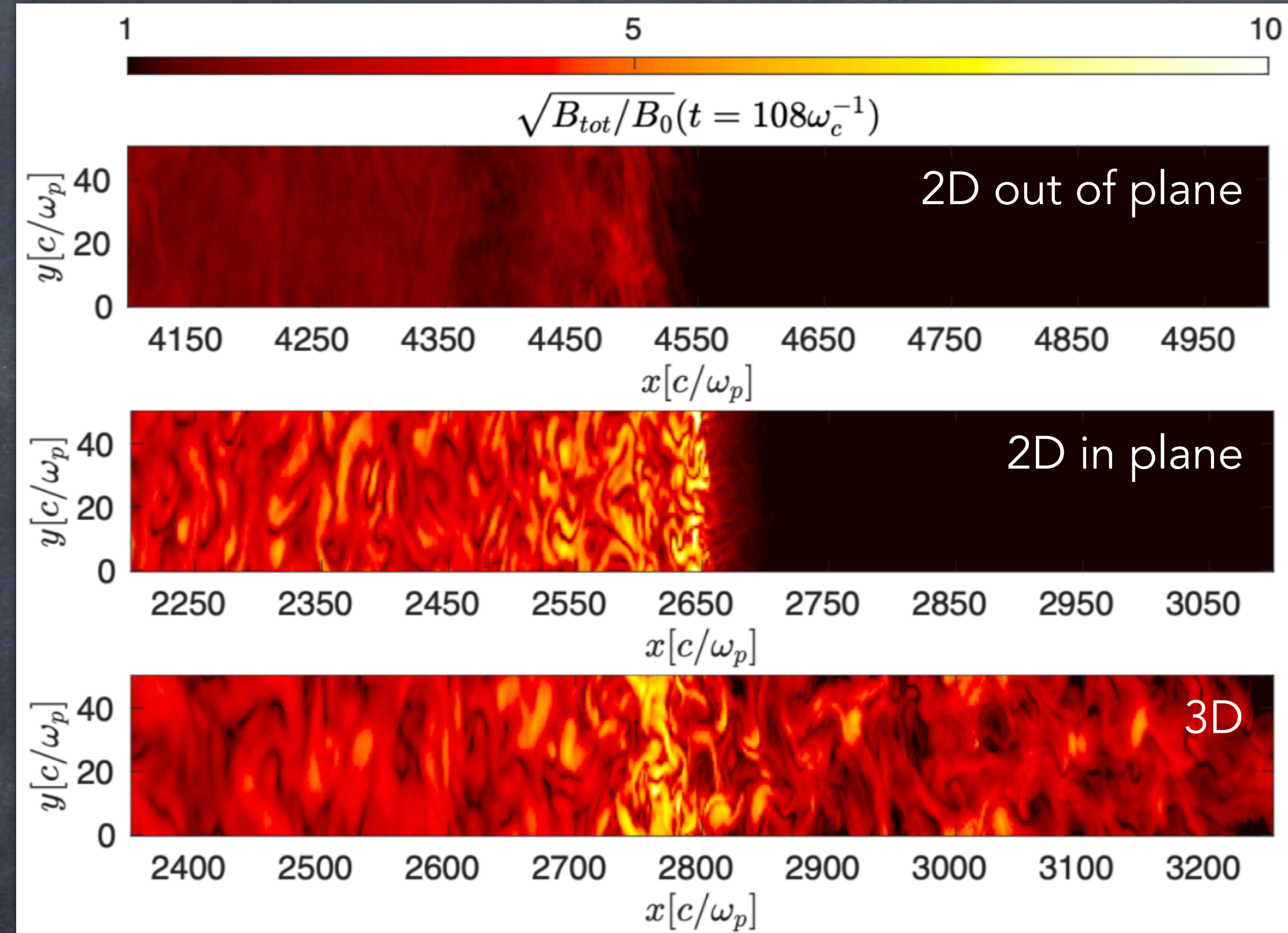
Caprioli & Spitkovsky14a,b

- Sironi+11 found $M_A^* \gtrsim 30$ for PIC relativistic shocks



Oblique Shocks: B-Field Amplification

2D/3D simulations of a shock with $M_A = 100$, $\theta_{Bn} = 80^\circ$ (Orusa & Caprioli 2023)



- Magnetic field generation:
 - 1D: simple compression (MHD)
 - 2D out-of-plane B_0 : \sim compression
 - 2D in-plane B_0 : $\delta B/B_0 \lesssim 40$ at the shock
 - 3D: $\delta B/B_0 \lesssim 40$ at the shock, but also $\delta B/B_0 \gg 1$ upstream
- Dimensionality matters! Why?
 - Turbulence is different in 3D...

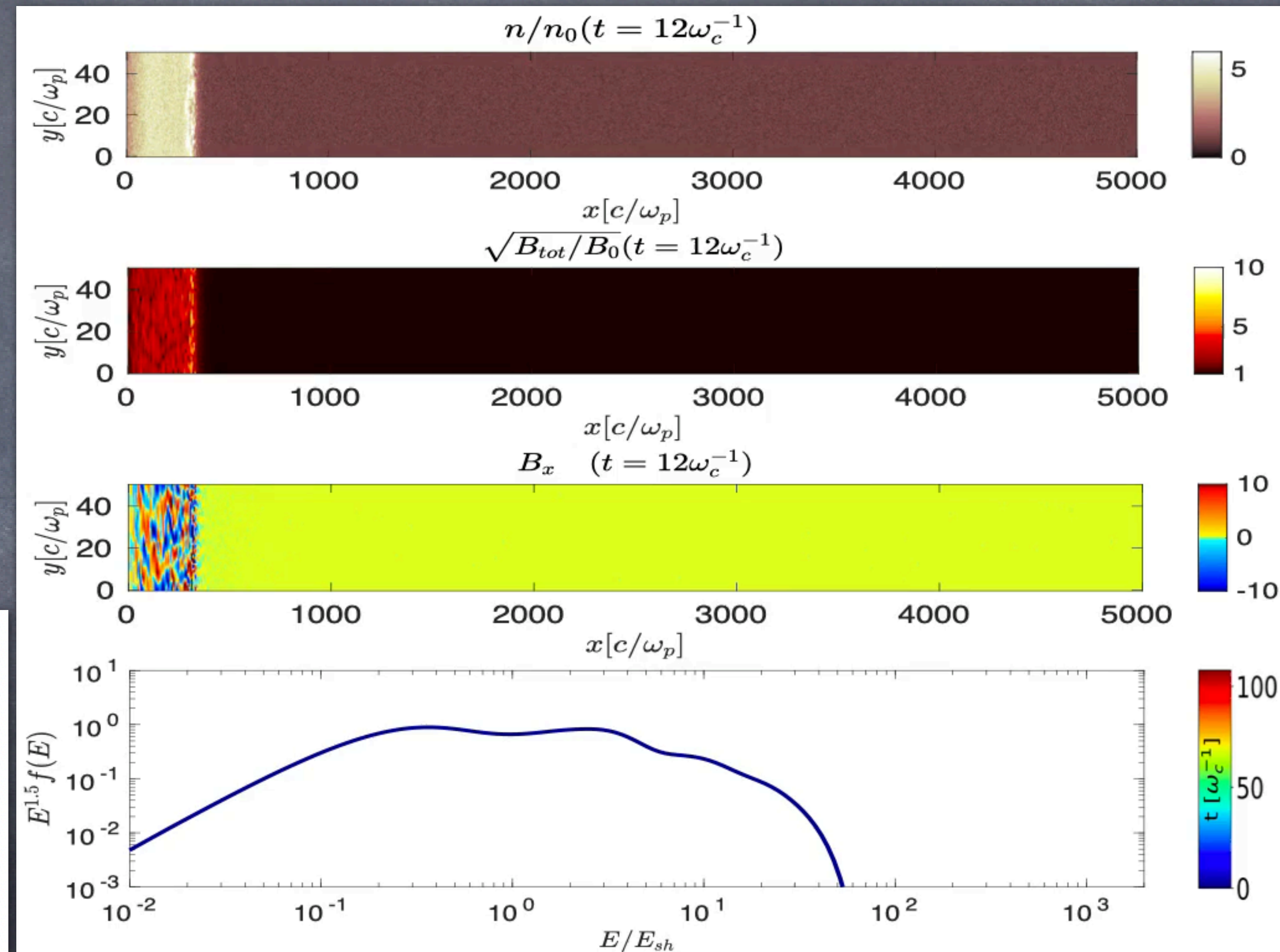
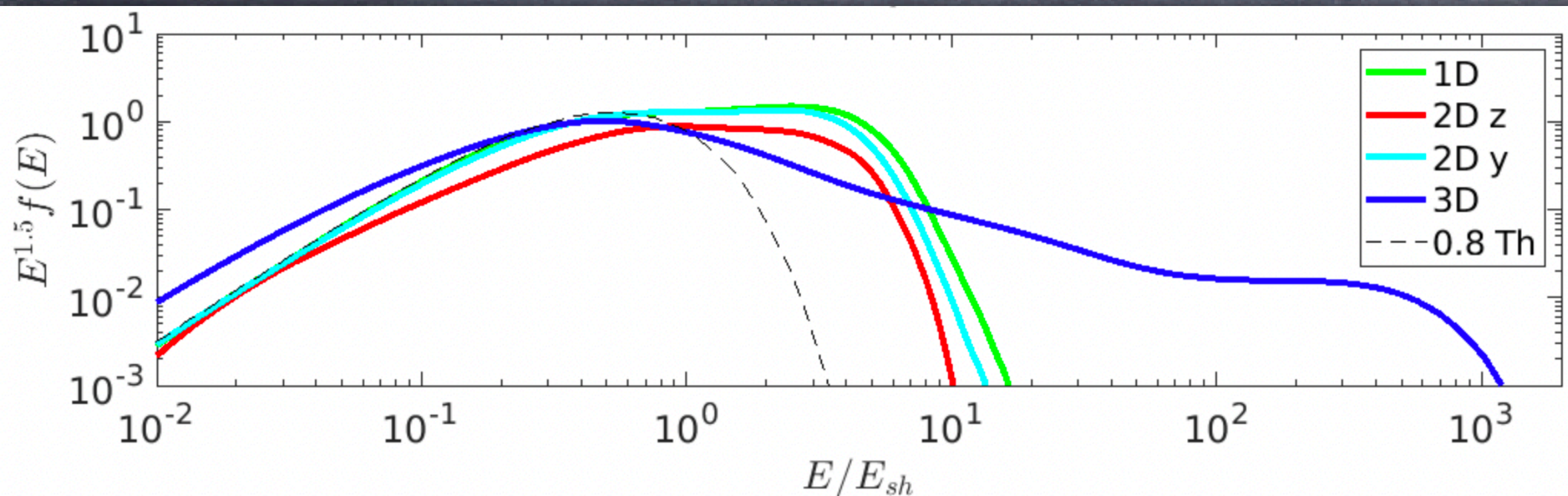
Oblique Shocks: Ion Acceleration

- Self-generated turbulence solves the injection problem!

- 3D geometry unlocks **cross-field diffusion** / B-field line wandering

- Supra-thermal ions can **diffuse back** from downstream

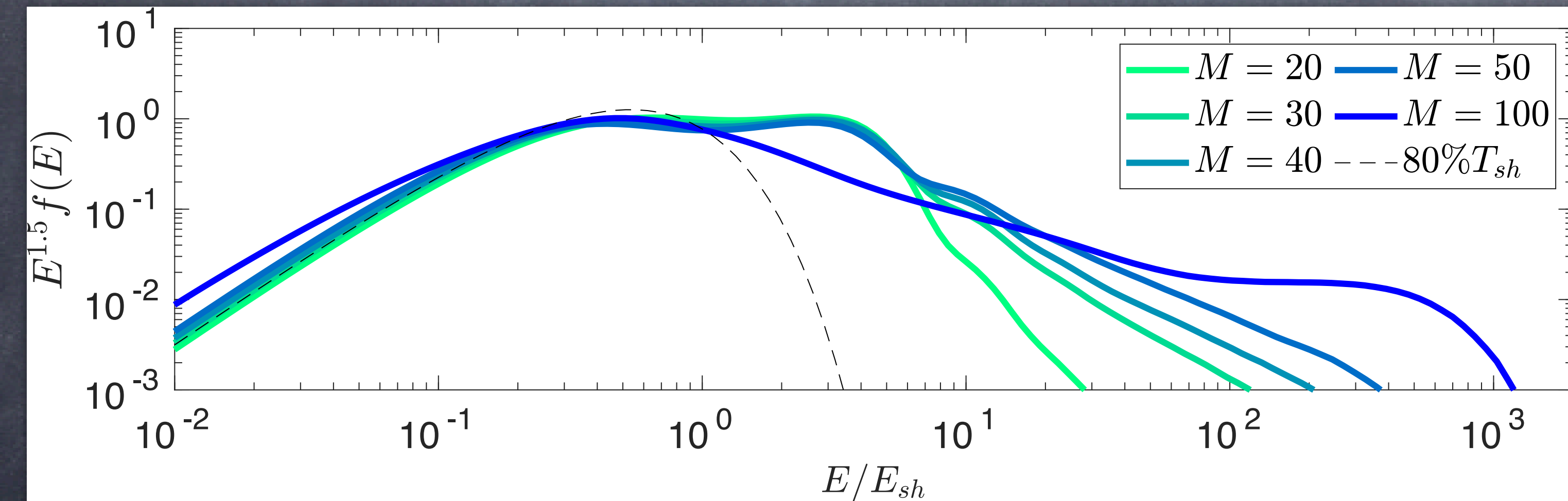
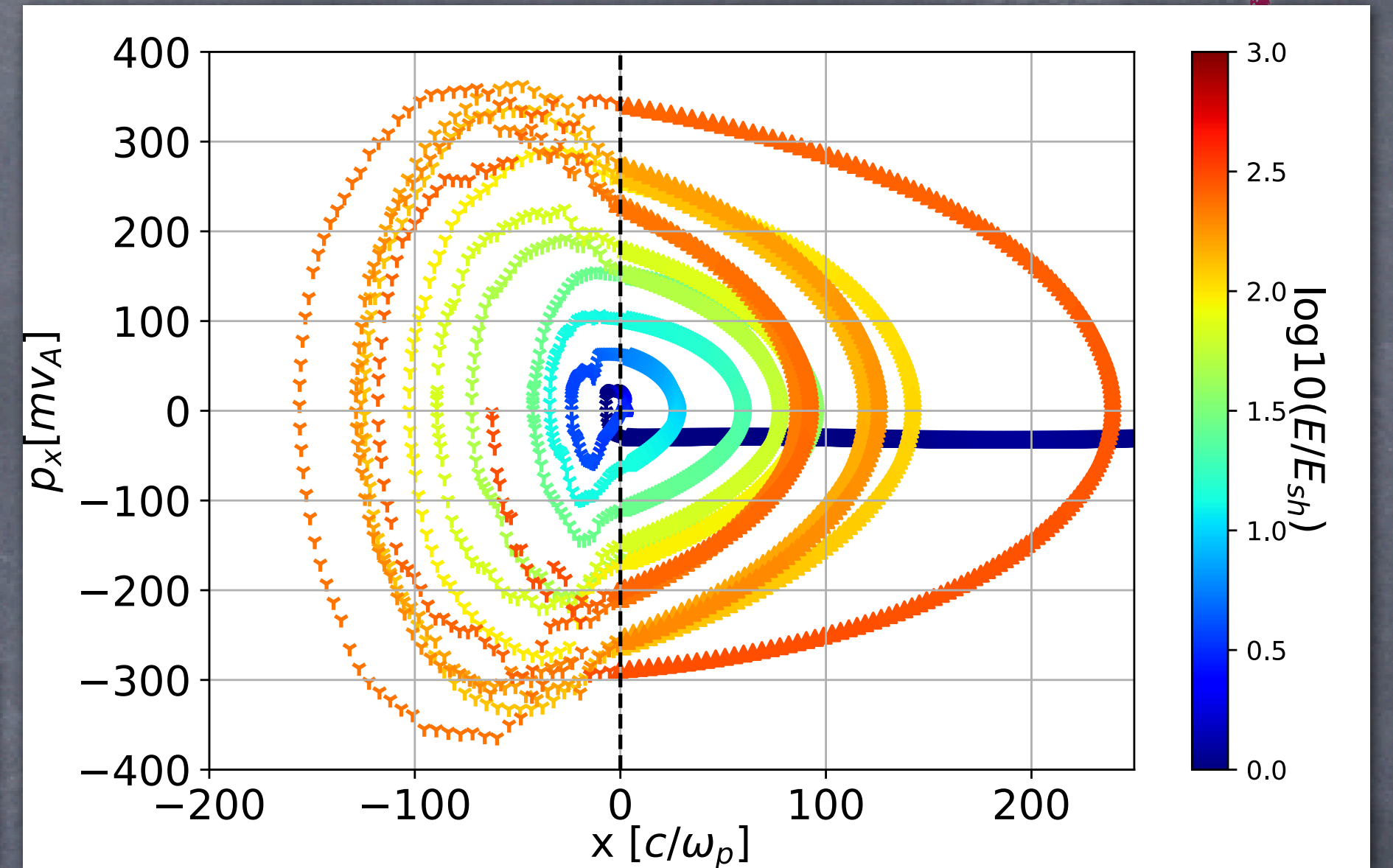
- and develop a **non-thermal tail**



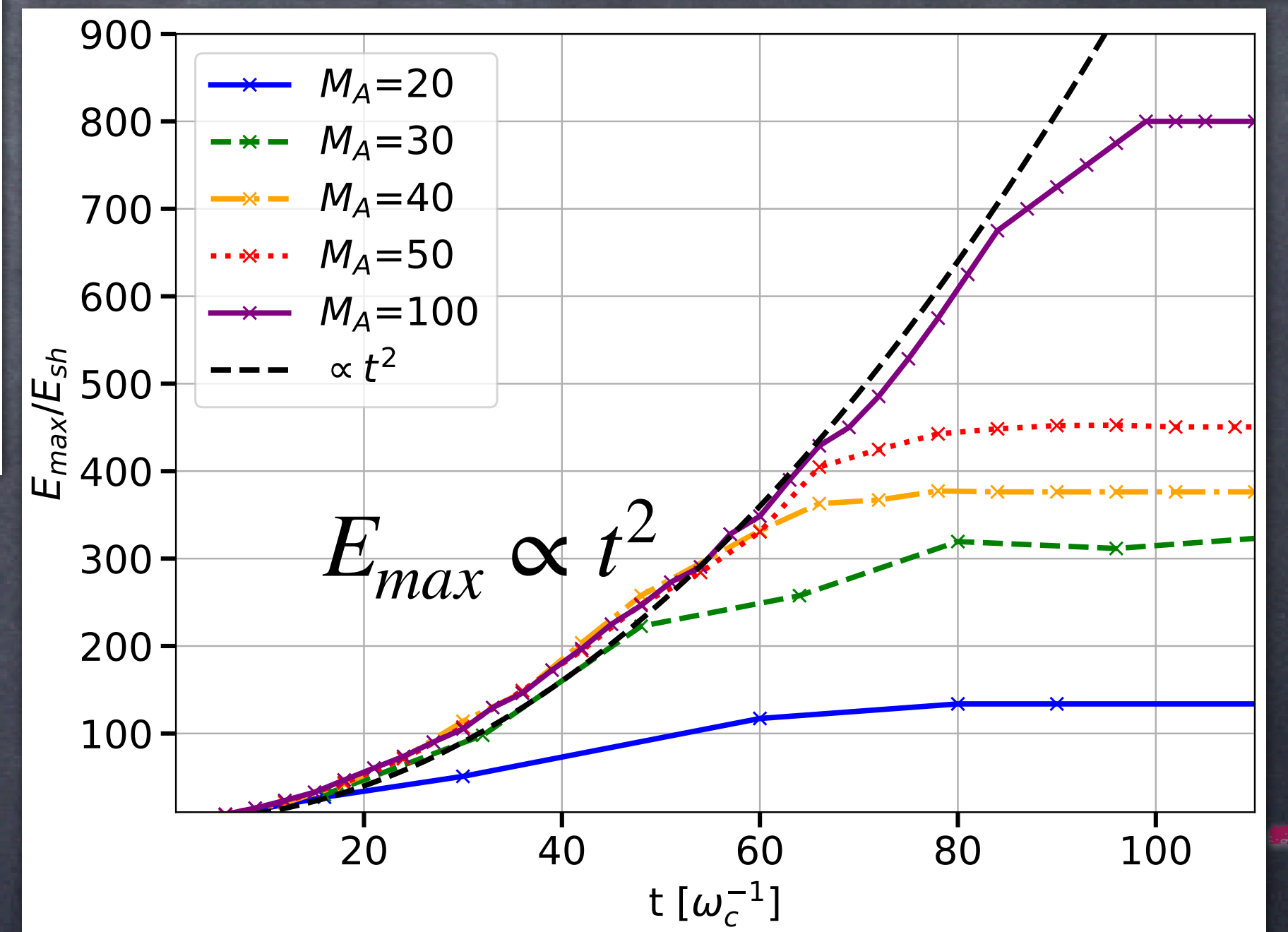
Oblique Shocks: Shock Drift Acceleration



- Particle tracking reveals that ions gain energy via **shock drift acceleration**
- Acceleration efficient ($\gtrsim 15\%$) and very fast!

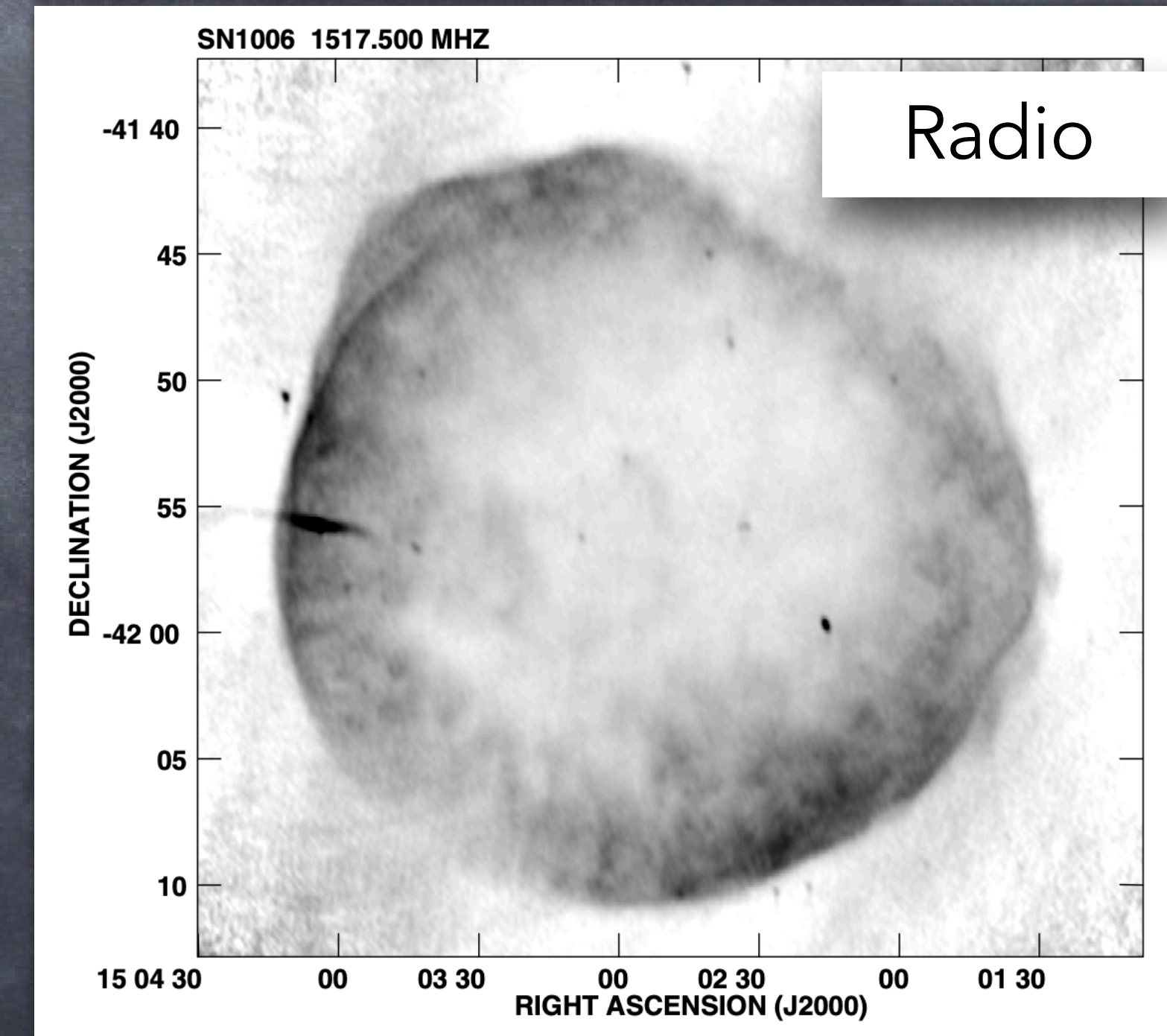


- There is maximum energy achievable via SDA
- Slope** and **maximum energy** depend on M_A



Implications for SNRs (e.g., SN1006)

- Investigate **dependence on θ_{Bn} and M_A** (w. Orusa, Simon, in prog.)
- Preliminary:** allows **injection also for oblique shocks** with $45^\circ < \theta_{Bn} < 65^\circ$
 - consistent with the compressions inferred in SN1006 (Giuffrida+21)
- SDA energy gain is limited ($\propto M_A$), then ions **escape** upstream
- They have hard time **driving Bell instability, so no DSA**
 - See **E. Simon's** talk tomorrow
- SDA is very fast:** $E_{max} \propto t^2$
 - A $v_{sh} \simeq 3,000$ km/s can make \sim GeV particles in $\lesssim 1$ day
 - Explains **azimuthally symmetric radio emission** from SN1006
- But is intrinsically limited to relatively small $E_{max}(\theta_{Bn})$
 - Explains **lack of X-ray synch and TeV emission**



TAKE-AWAY MESSAGES

- **Particle acceleration** is generally efficient in SNR shocks
- **Q-parallel**: DSA is efficient ($\gtrsim 10\%$)
 - Efficient **B amplification** via Bell's instability (Caprioli-Spitkovsky14a,b,c)
 - **Slope** not universal! Steeper than E^{-2} , depends on B (Caprioli+20, Diesing & Caprioli 21)
 - E_{max} determined by the time it takes to grow B (Simon's talk, in prog.)
- **Q-perpendicular**: SDA efficient if $M_A \gtrsim 30$; no injection problem (Orusa & Caprioli+20)
 - **Slope** not universal! Steep, but $\rightarrow E^{-2}$ for $M_A \gtrsim 100$
 - Generally **limited** to $E_{max} \lesssim 10$ GeV (Orusa & Caprioli, in prog.)
- Use shock acceleration theory to interpret **SNR multi-wavelength** emission!

